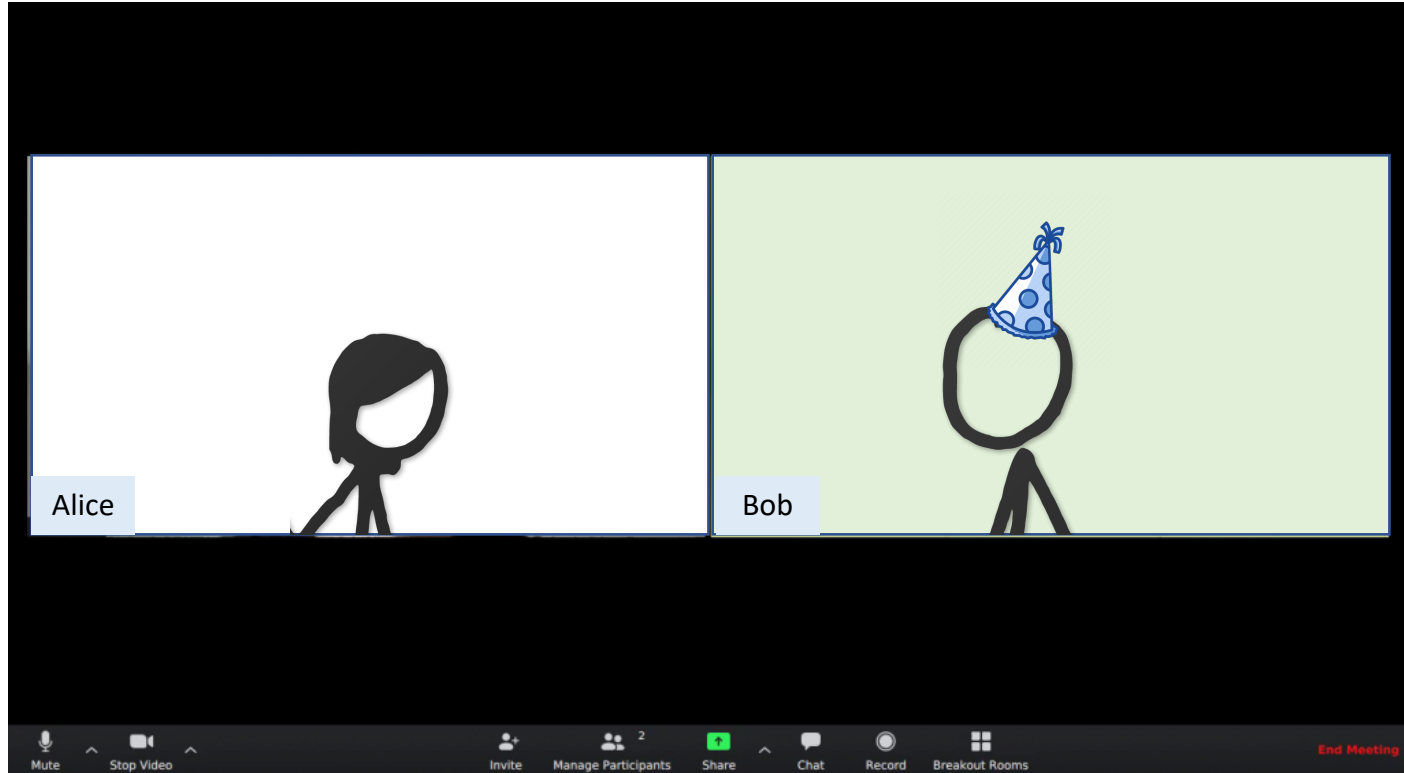


Crash Course in Quantum Computing

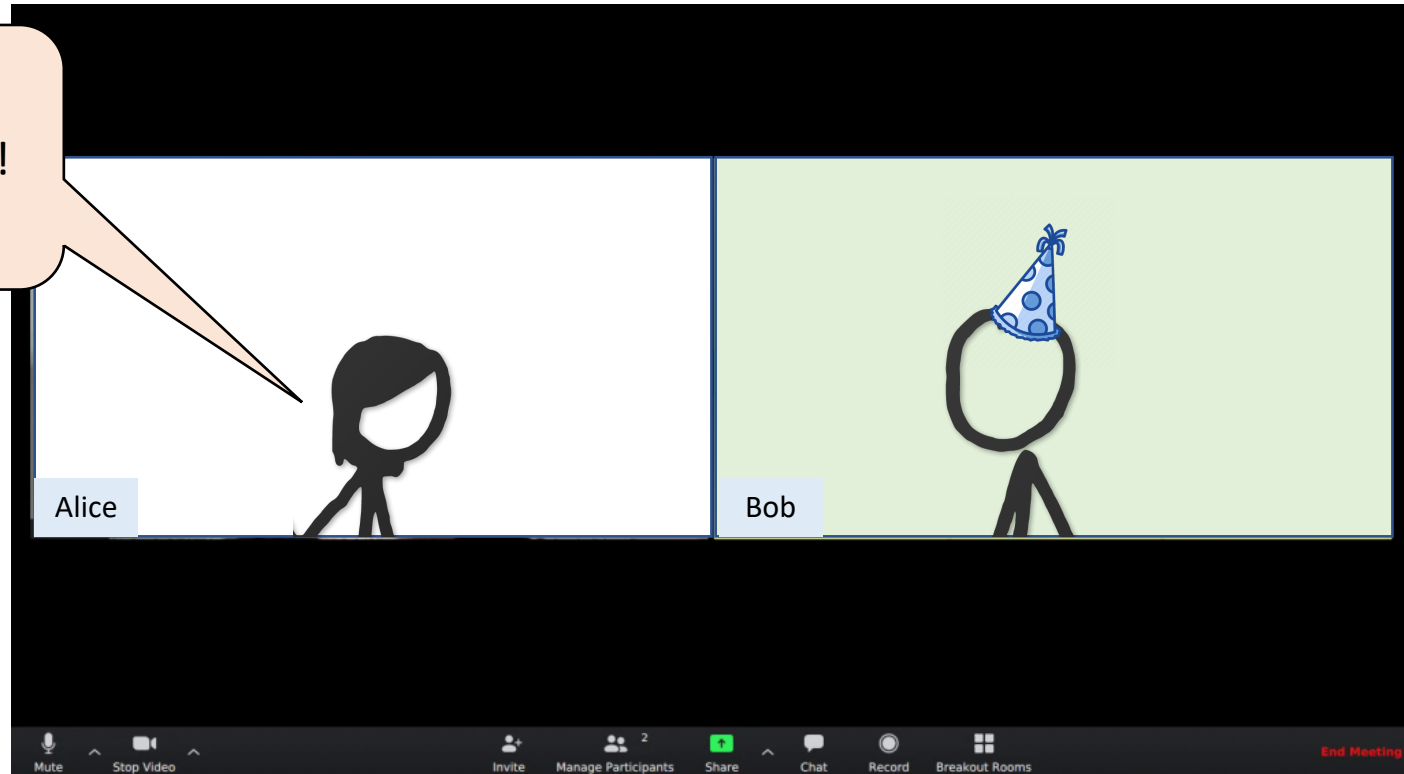
Hour 2: Quantum Computation

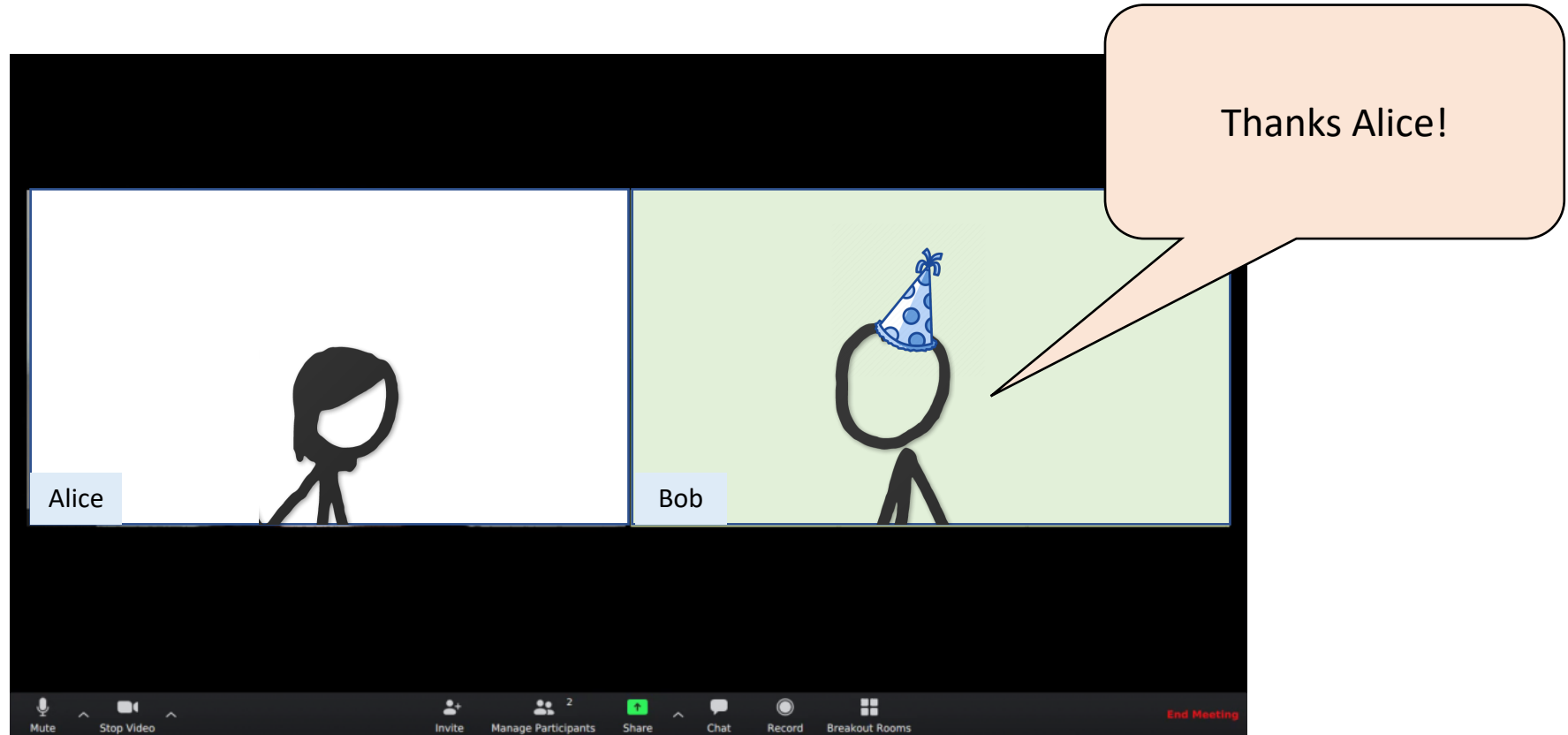
BIU Winter School on Cryptography 2021

Lecturer: Henry Yuen

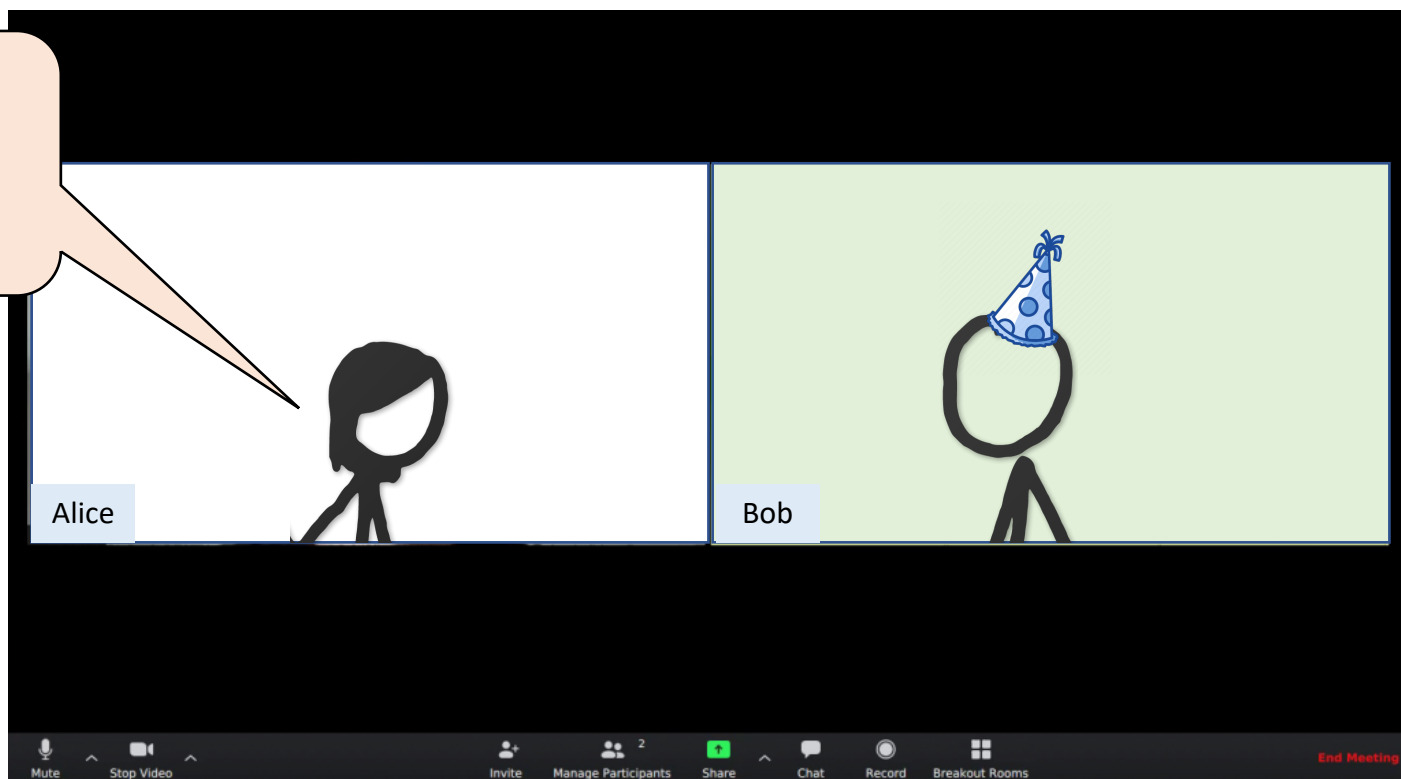


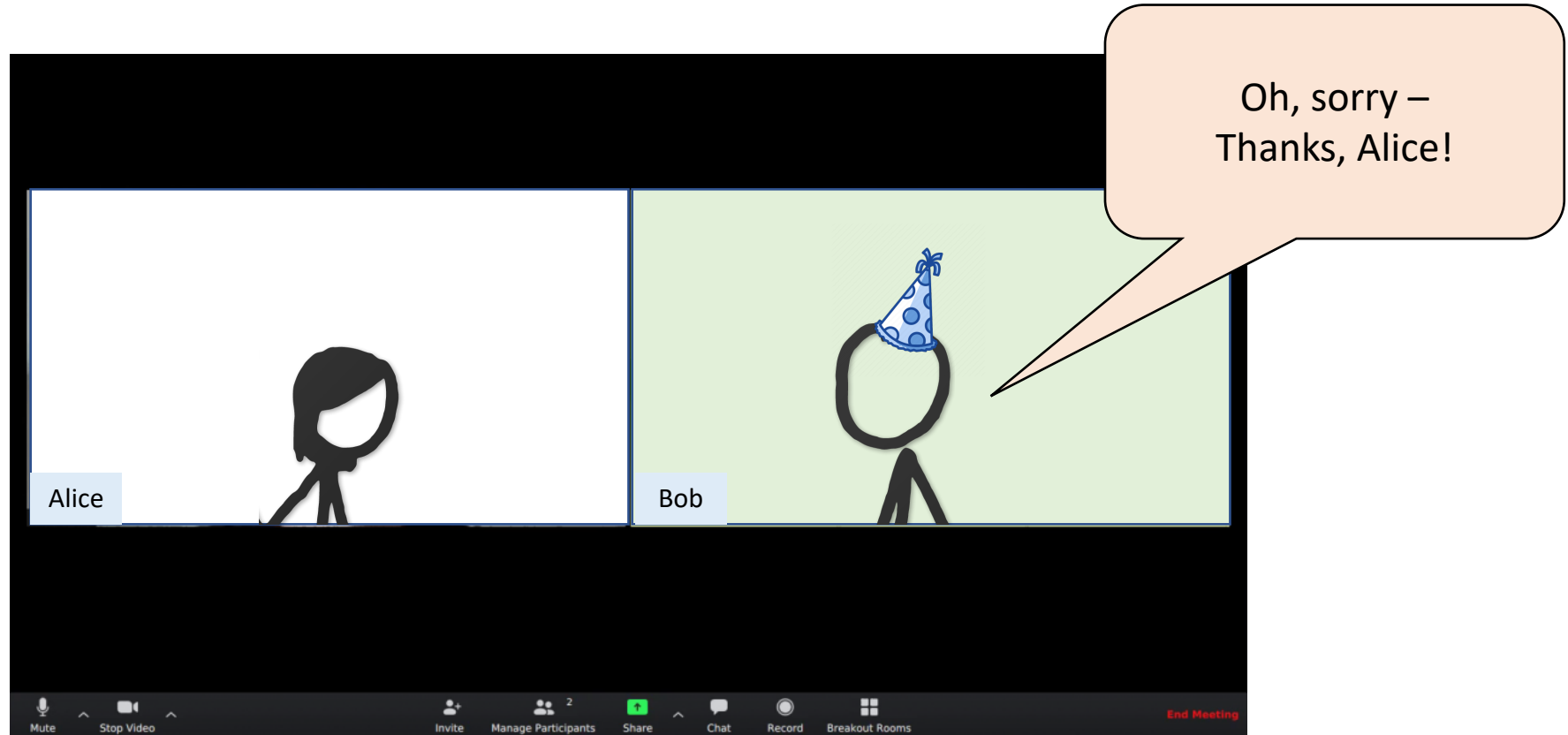
Happy Birthday, Bob!



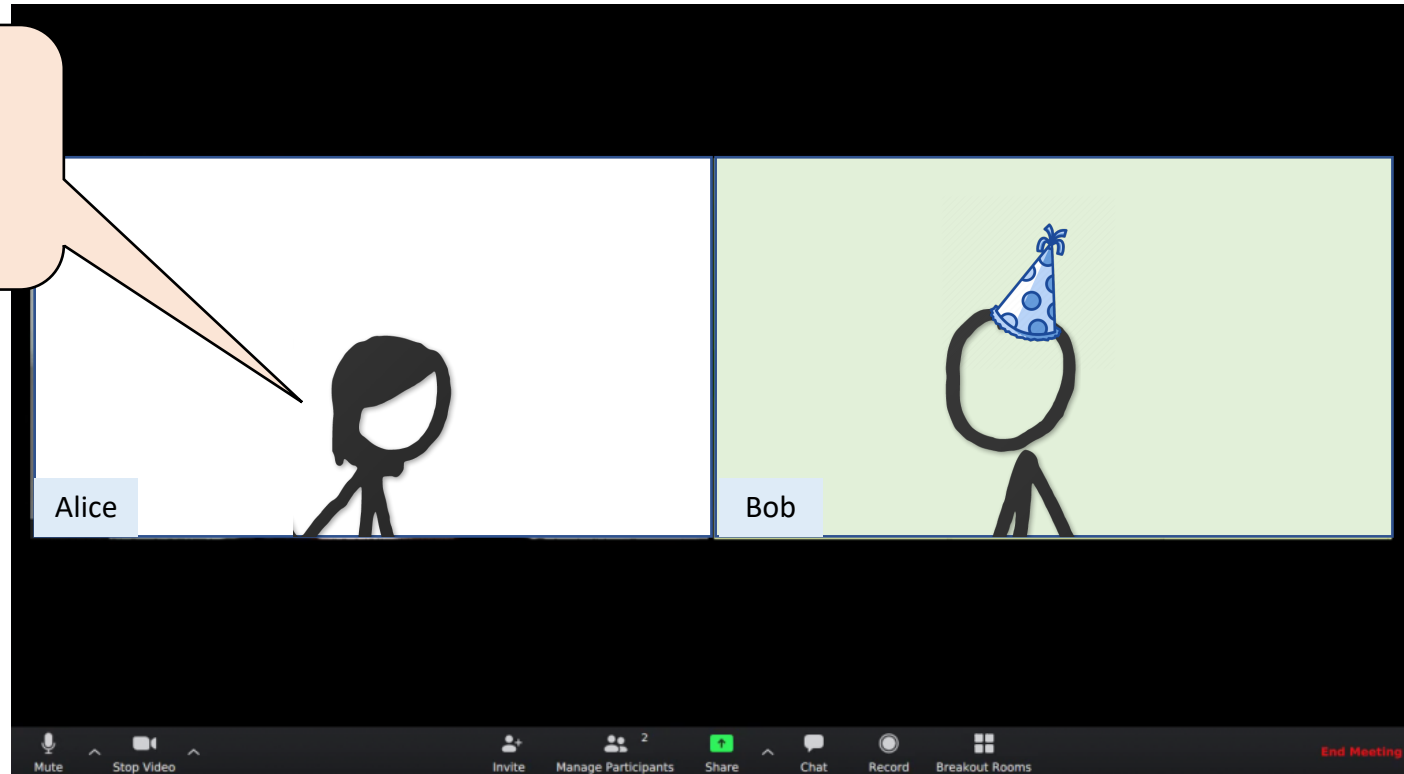


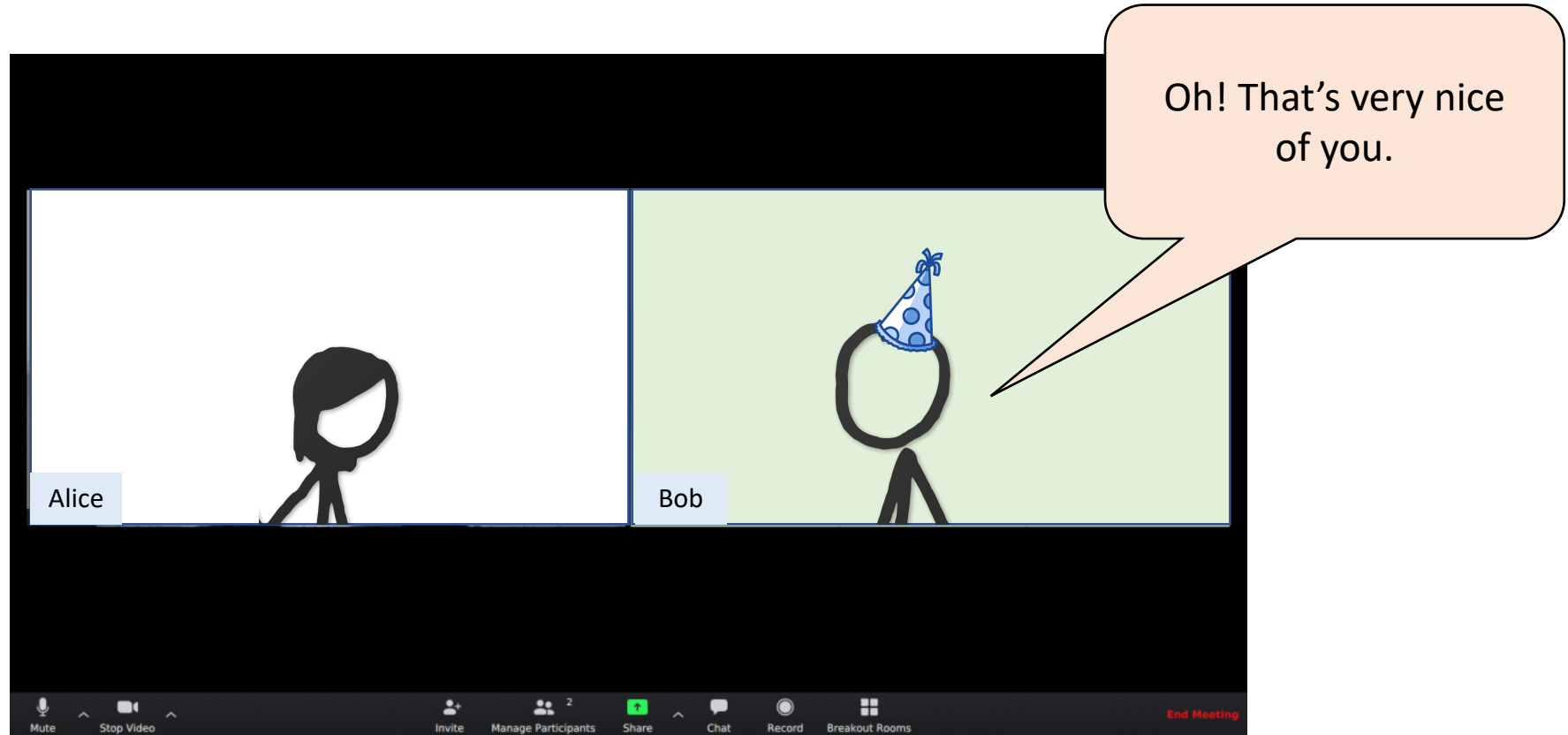
You're muted.



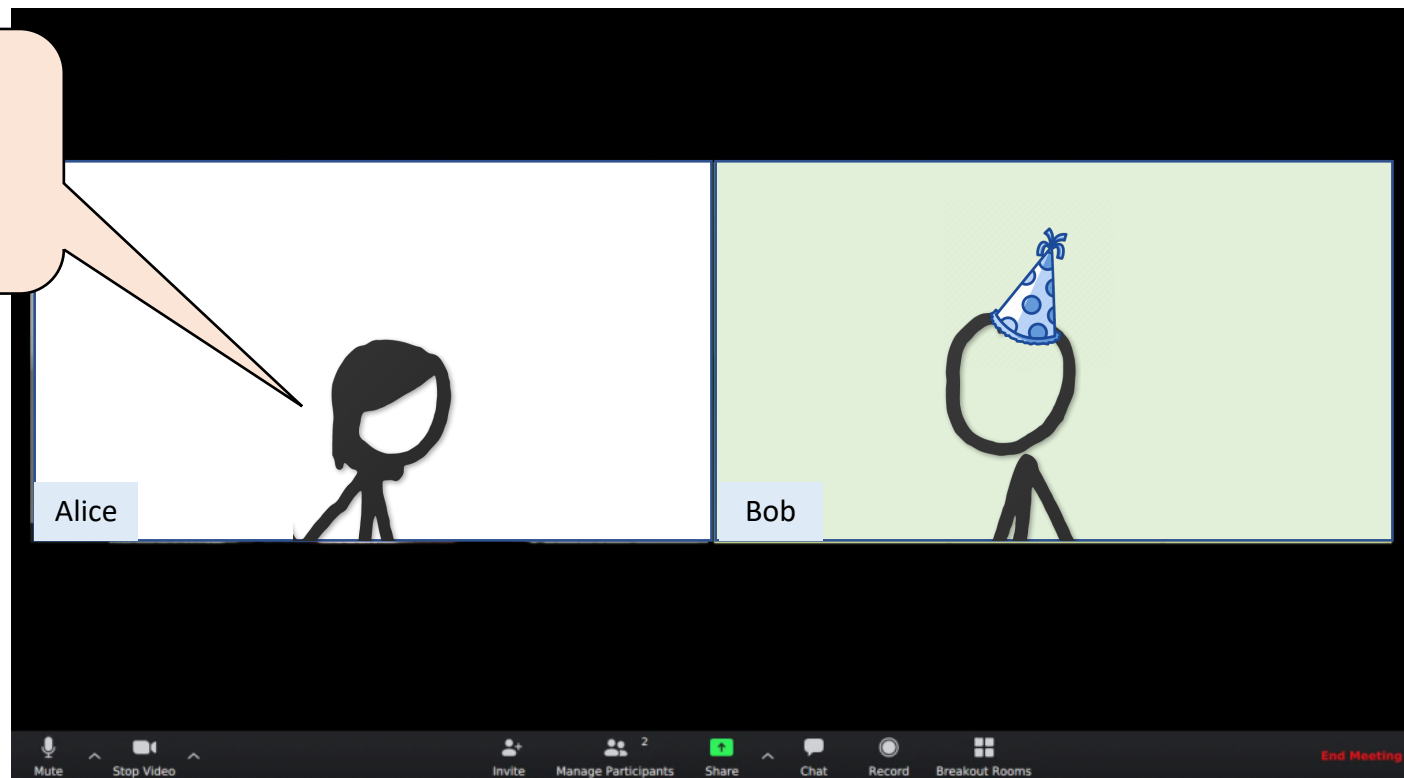


I have a gift for you.

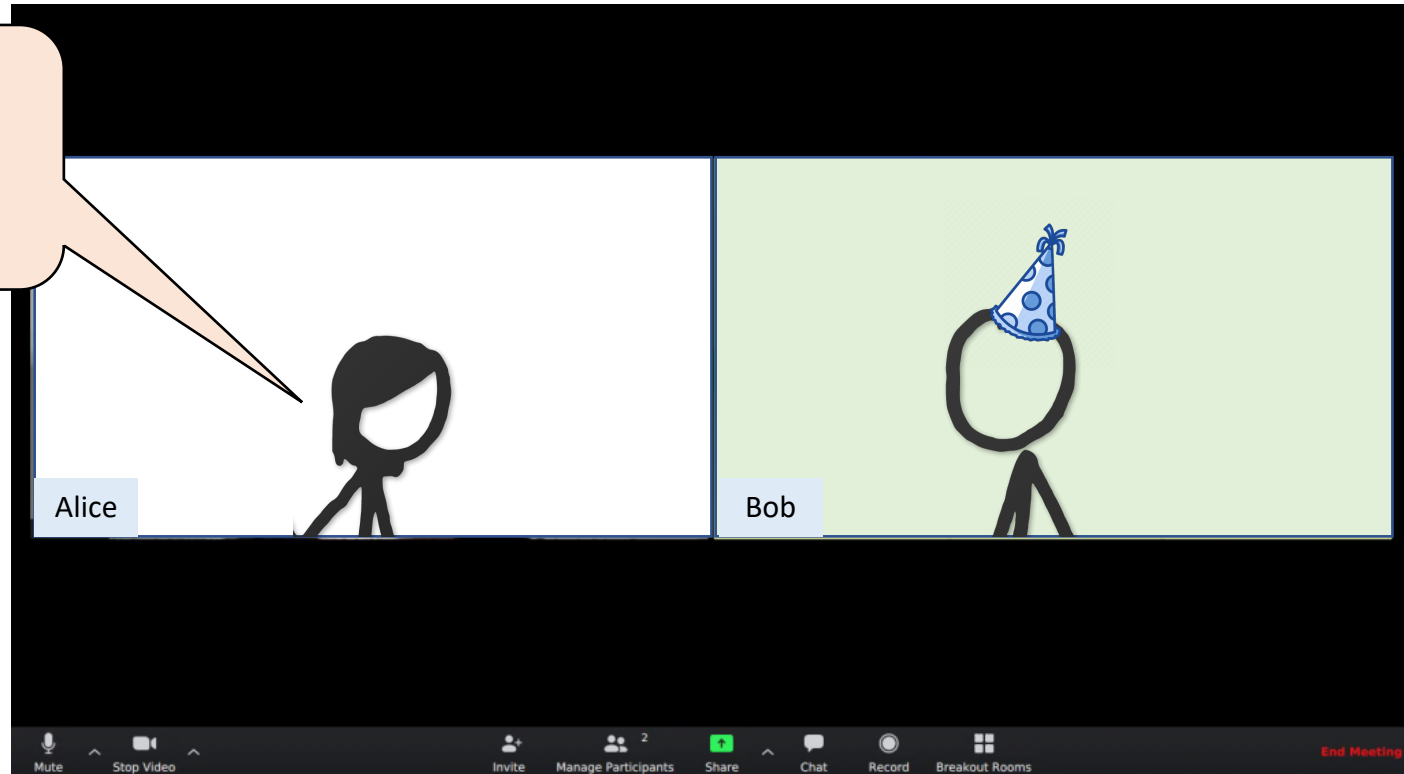


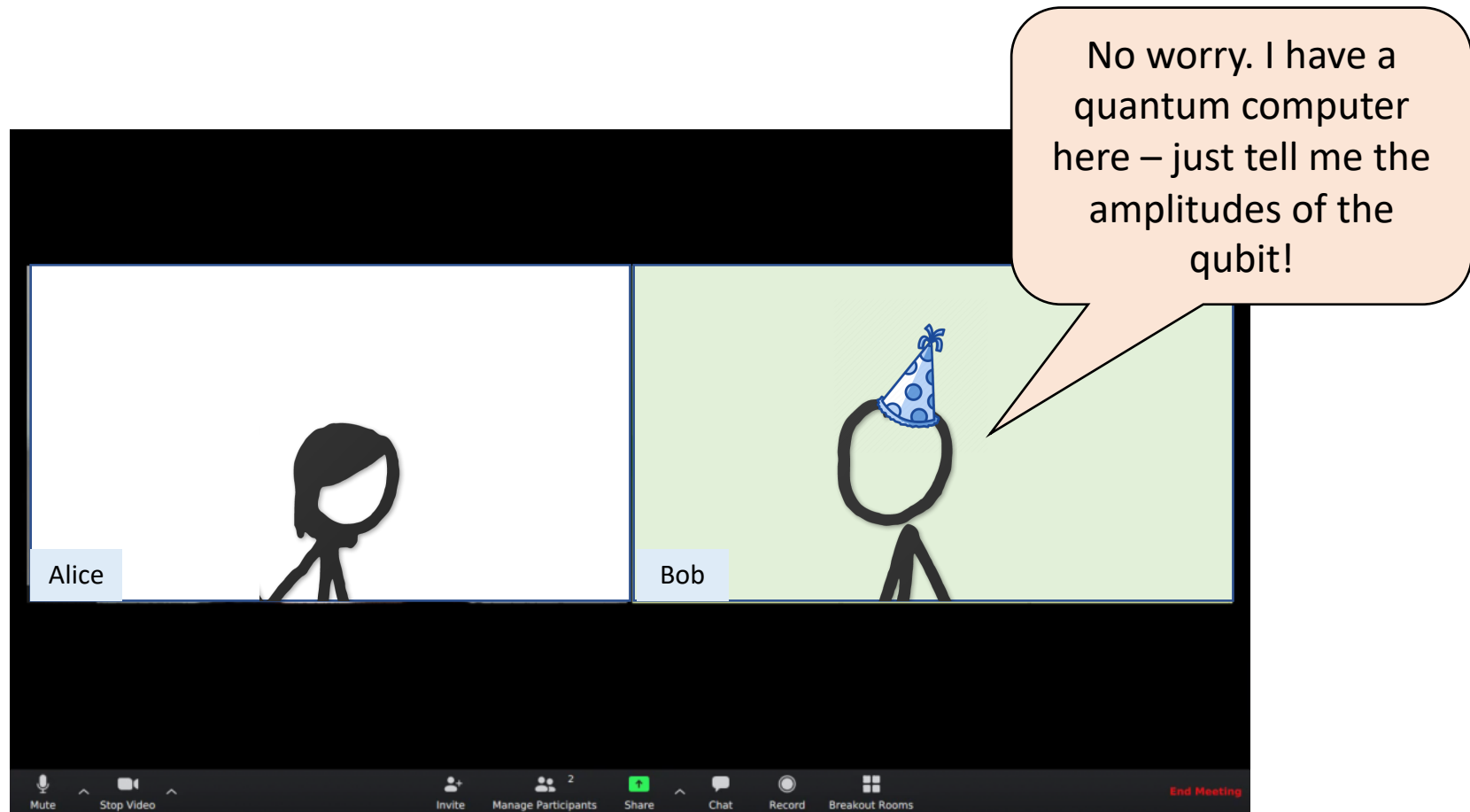


There's a problem,
though.

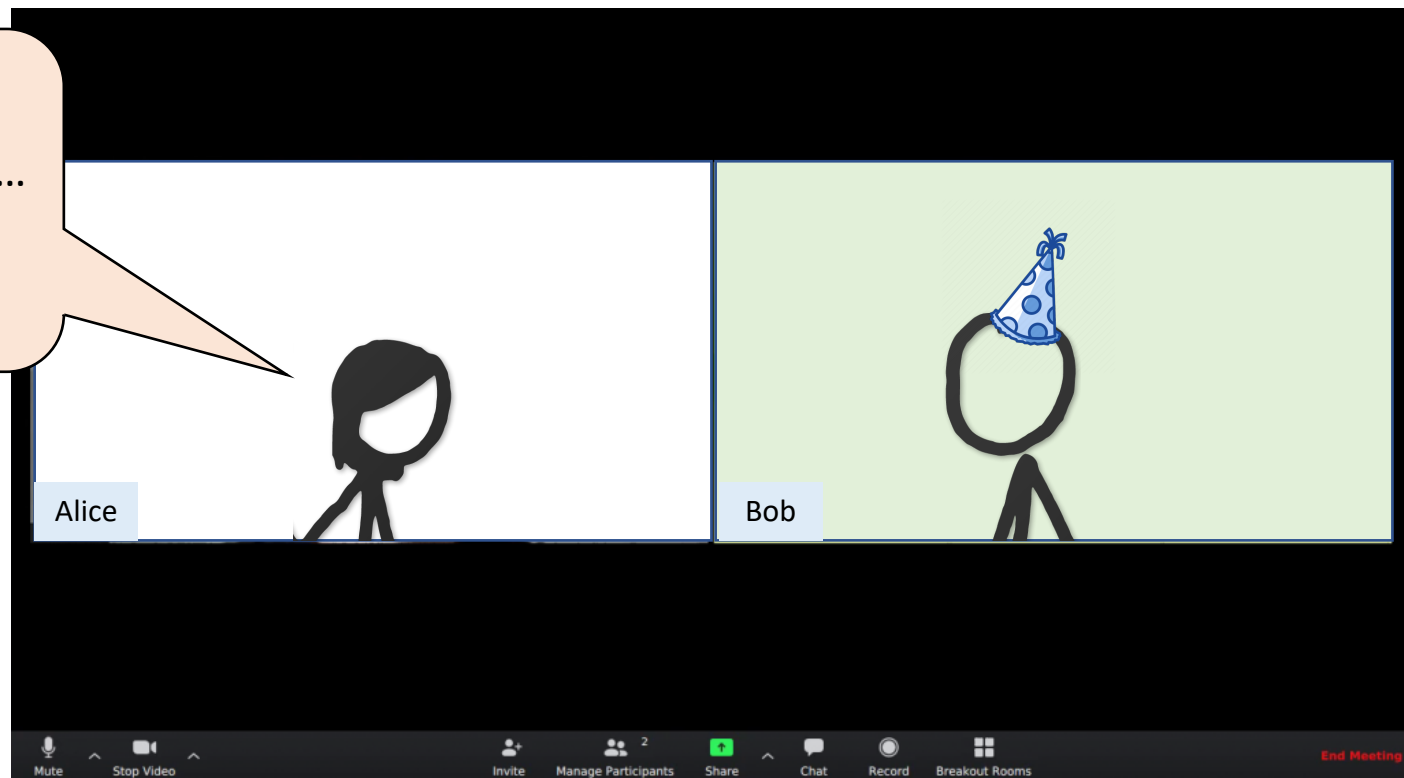


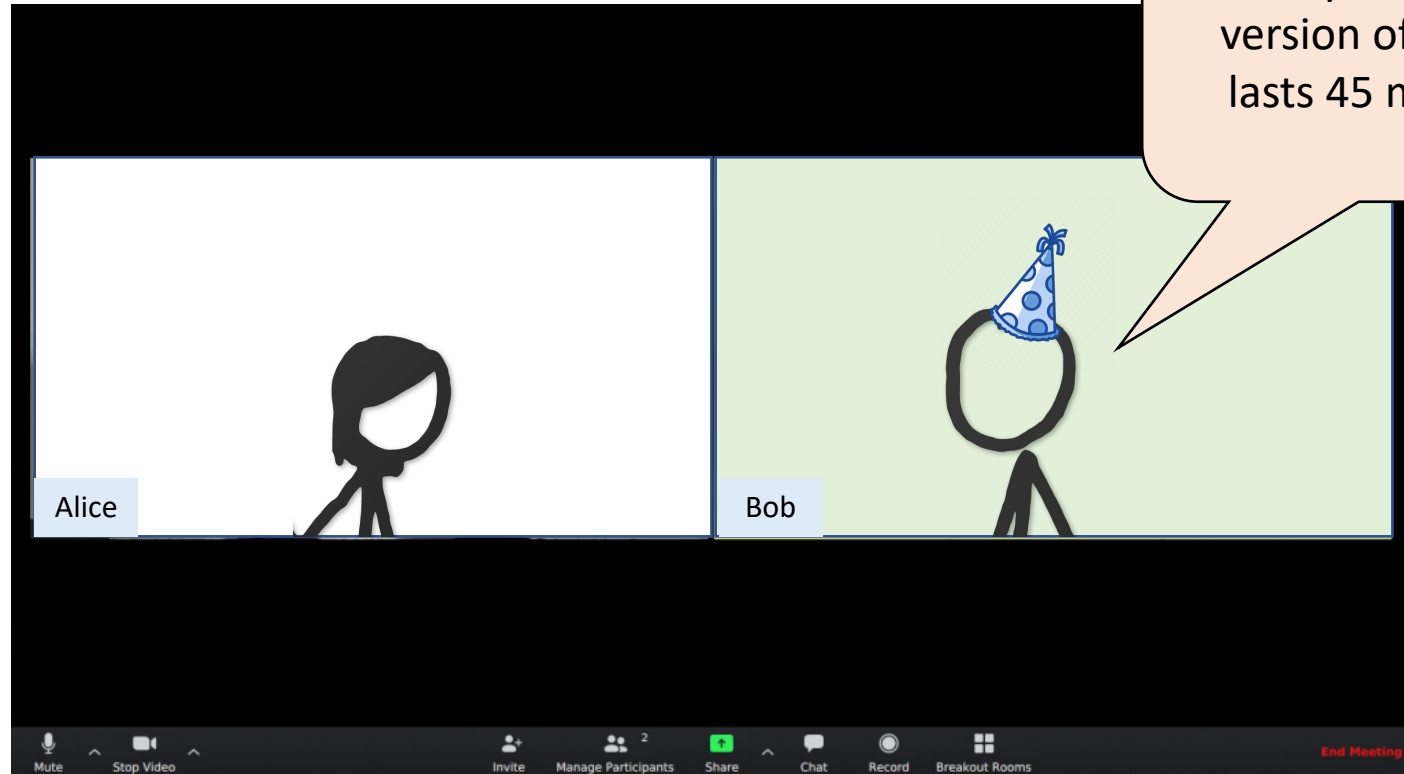
It's a qubit.

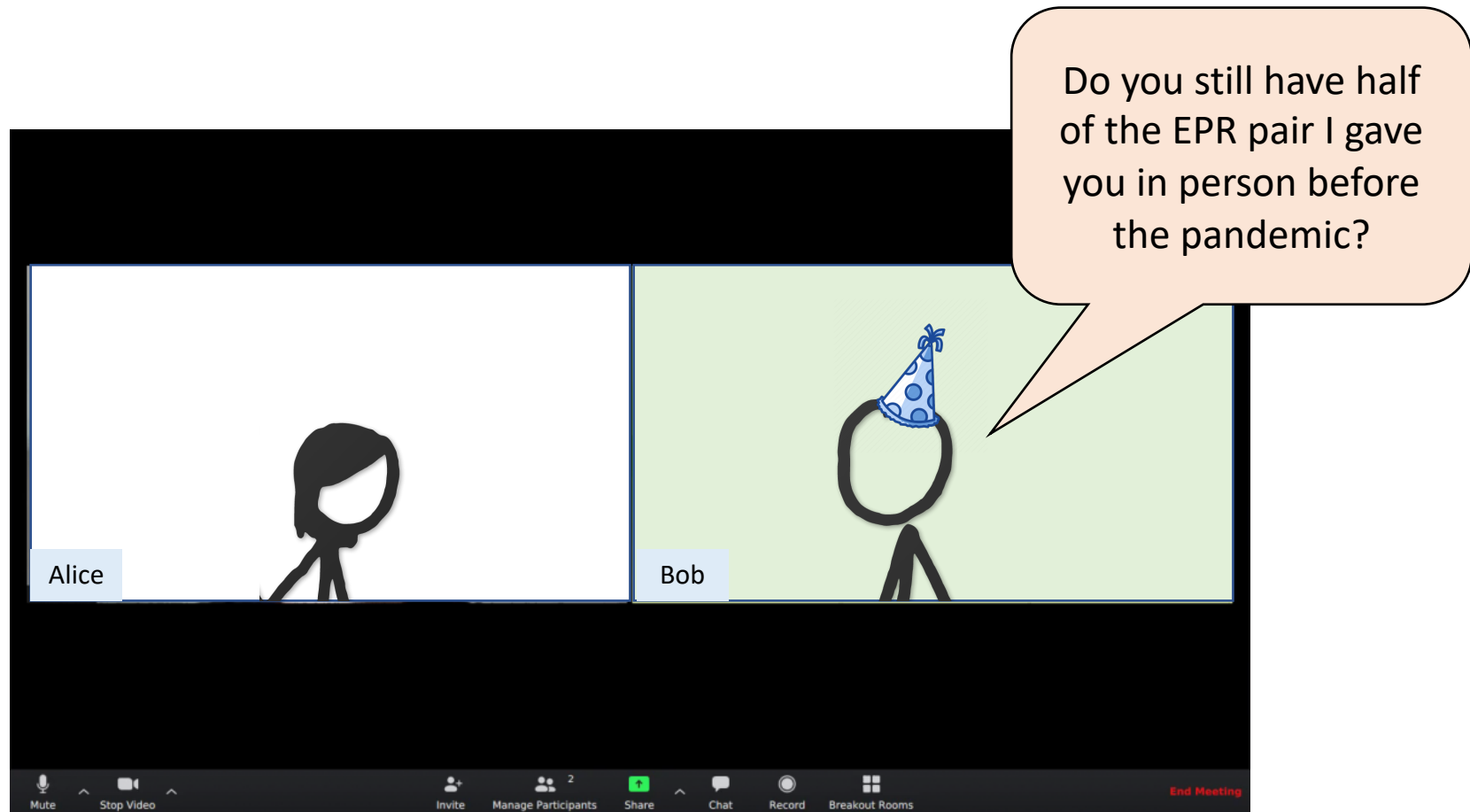




The amplitudes are
transcendental numbers...
I don't think our Zoom
call can last that long.

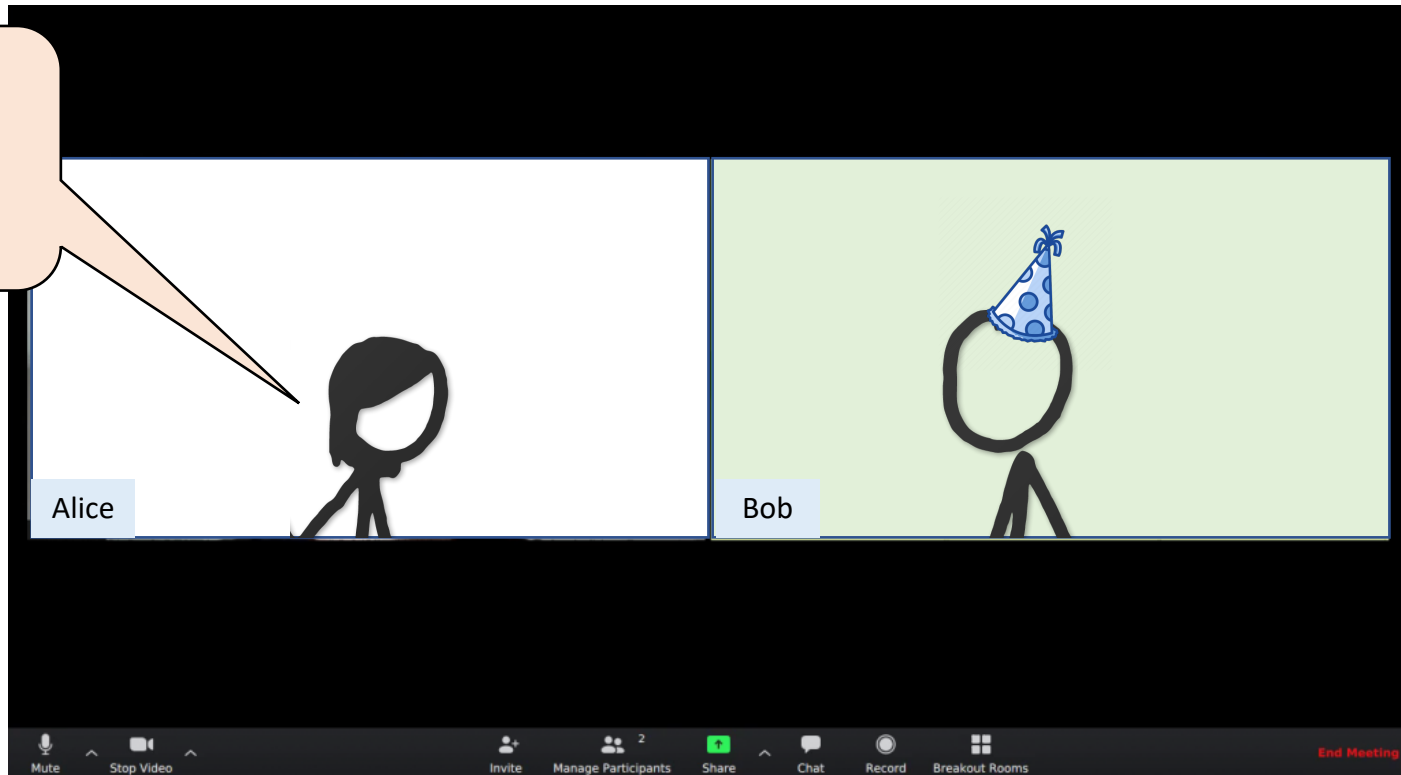




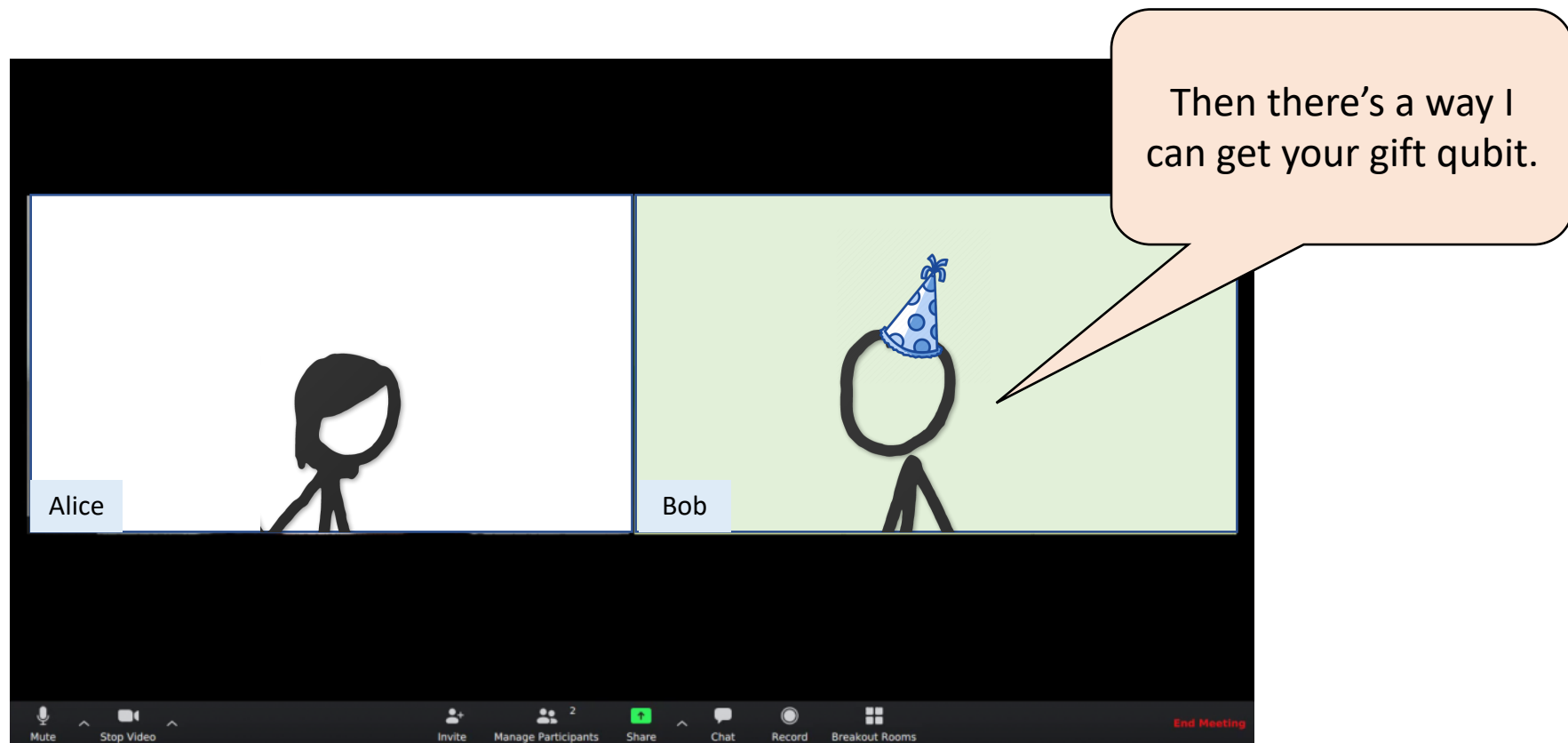


$$|EPR\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

I do, in fact!
That seems like
so long ago.

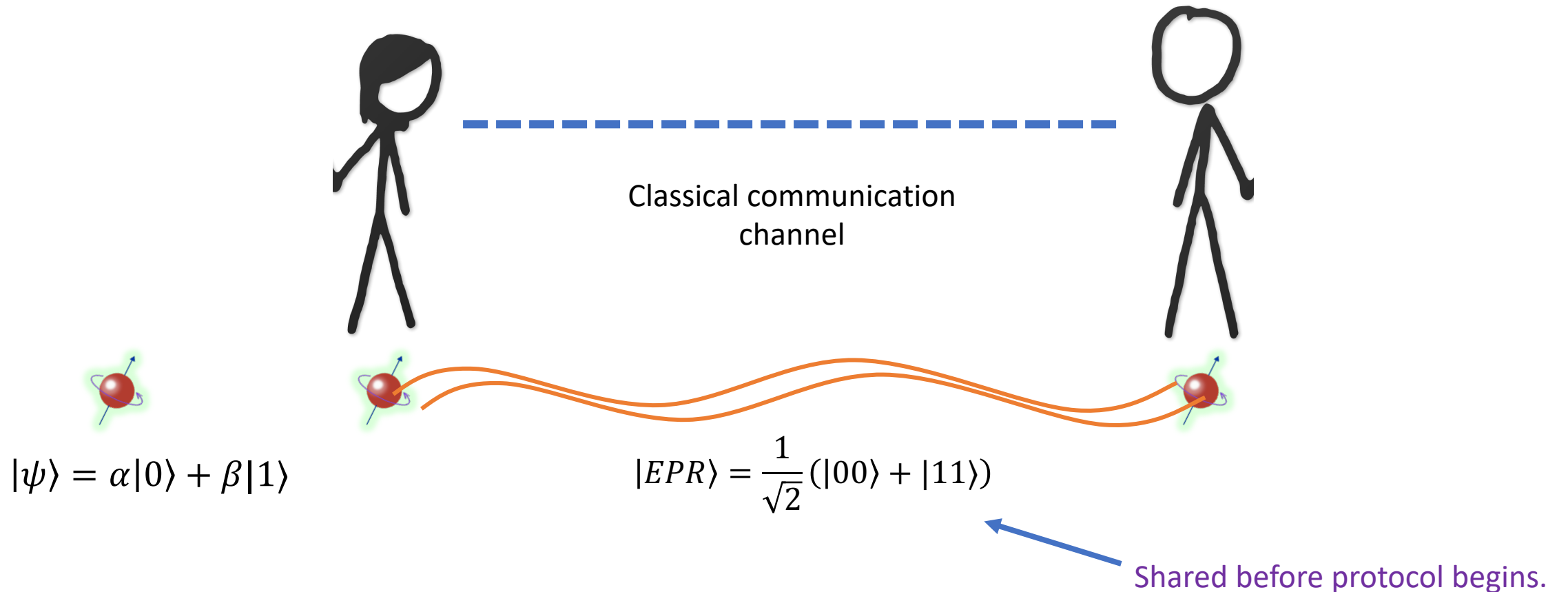


$$|EPR\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$



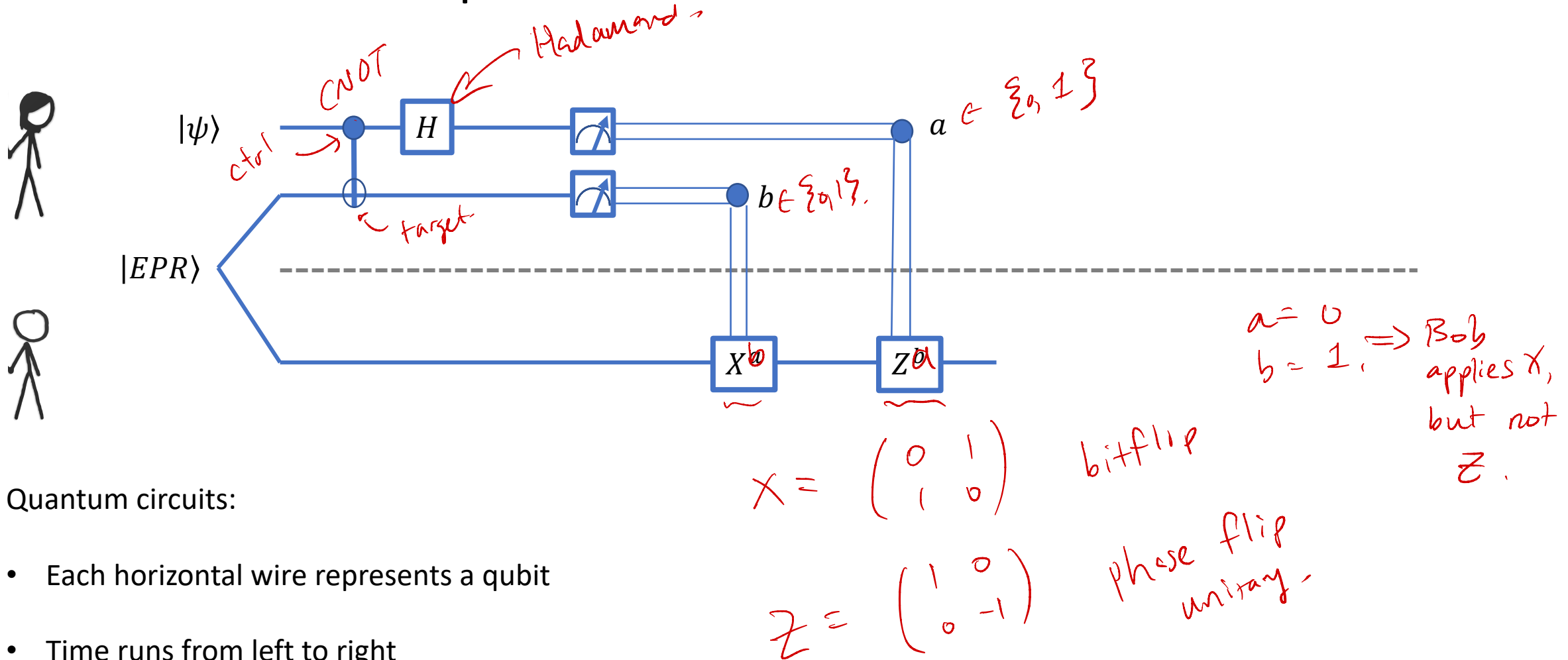
$$|EPR\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

Quantum Teleportation



Quantum teleportation allows Alice to send $|\psi\rangle$ to Bob using preshared entanglement and classical communication.

Quantum Teleportation



Quantum circuits:

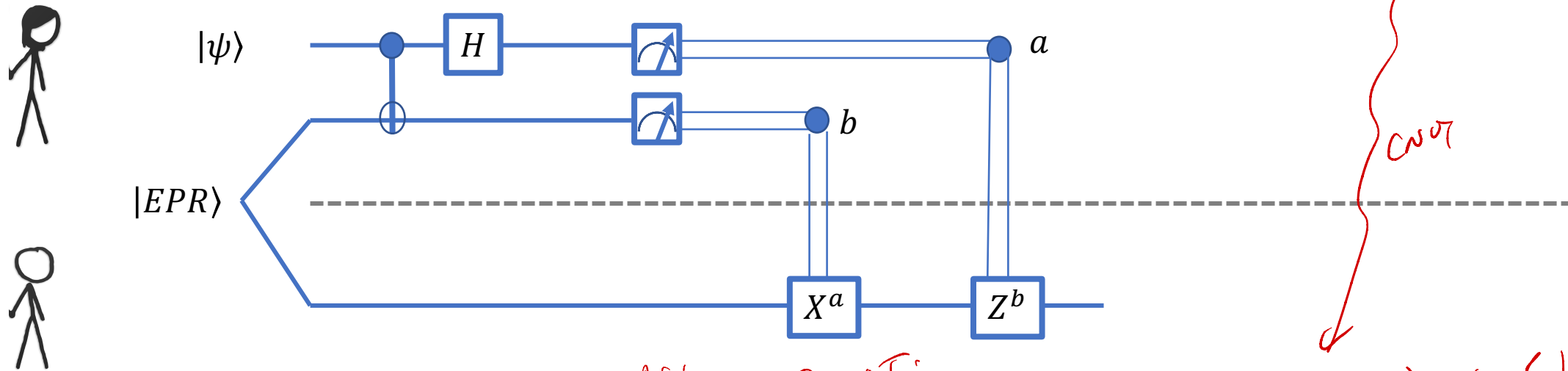
- Each horizontal wire represents a qubit
- Time runs from left to right
- Initial state of qubits is written on left hand side

Quantum Teleportation

Beginning:

$$|\psi\rangle \otimes |EPR\rangle = \frac{1}{\sqrt{2}} (\alpha|0\rangle + \beta|1\rangle) \otimes (|00\rangle + |11\rangle)$$

$$= \frac{1}{\sqrt{2}} (\alpha|000\rangle + \alpha|011\rangle + \beta|100\rangle + \beta|111\rangle)$$



Quantum circuits:

- Each horizontal wire represents a qubit
- Time runs from left to right
- Initial state of qubits is written on left hand side

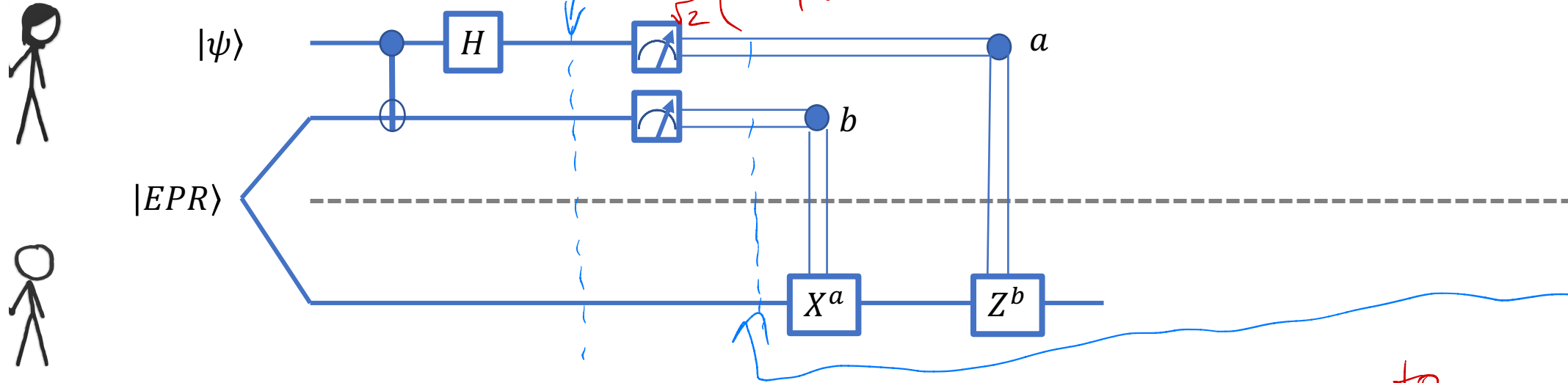
After CNOT:

$$\frac{1}{\sqrt{2}} (\alpha|000\rangle + \alpha|011\rangle + \beta|110\rangle + \beta|101\rangle)$$

After Hadamard:

$$\frac{1}{\sqrt{2}} (\alpha|+\rangle|00\rangle + \alpha|+\rangle|11\rangle + \beta|-\rangle|10\rangle + \beta|-\rangle|01\rangle)$$

Quantum Teleportation

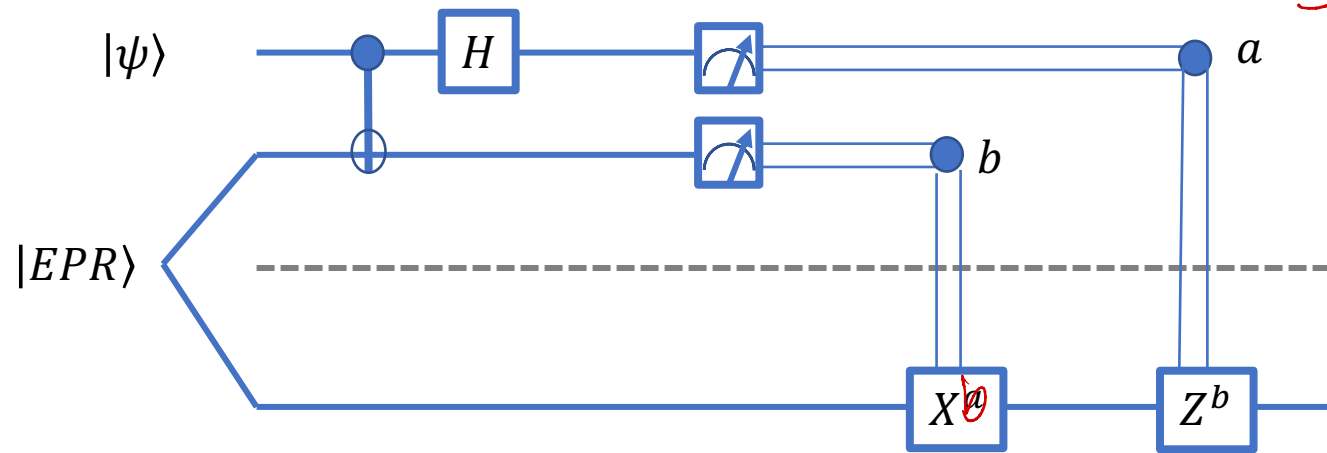


Claim: At the end of protocol, Bob has $|\psi\rangle$.

suppose $b=1$. state collapses to,

$$\begin{aligned}
 & \alpha |+-\rangle |11\rangle + \beta |-\rangle |10\rangle \\
 &= \frac{1}{\sqrt{2}} (\alpha |011\rangle + \beta |010\rangle + \alpha |111\rangle + \beta |110\rangle) \text{ Now measure first qubit.} \\
 & \text{suppose } a=0, \rightarrow \alpha |011\rangle + \beta |010\rangle = |0\rangle \otimes |1\rangle \otimes (\alpha |1\rangle + \beta |0\rangle)
 \end{aligned}$$

Quantum Teleportation



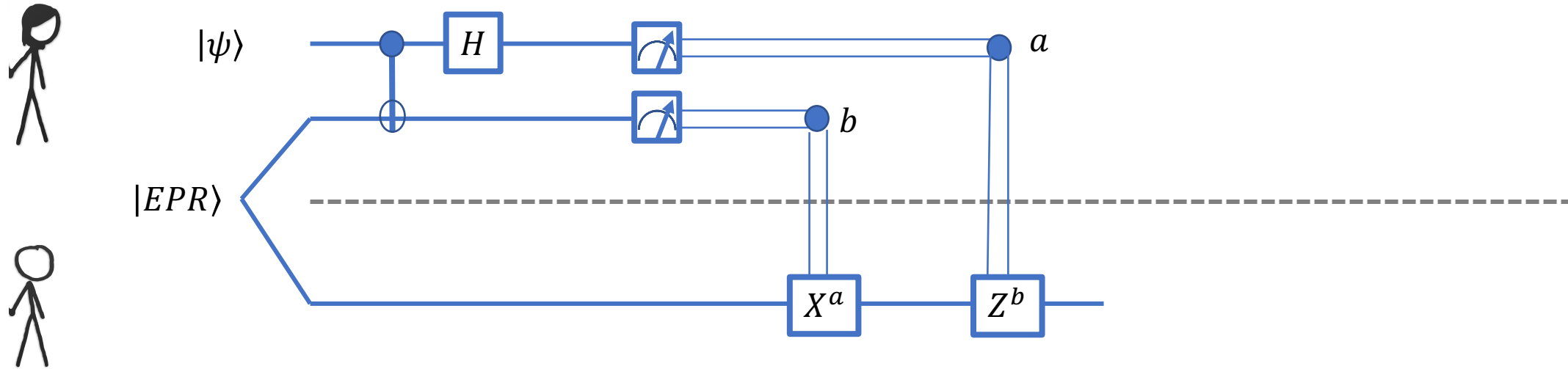
In the case $a=0, b=1$,
the state after measurement
is $|0, 1\rangle \otimes (\alpha|1\rangle + \beta|0\rangle)$.

→ Bob applies X to
his qubit -
→ Bob's qubit is in the
state

$$\alpha|0\rangle + \beta|1\rangle = |\psi\rangle.$$

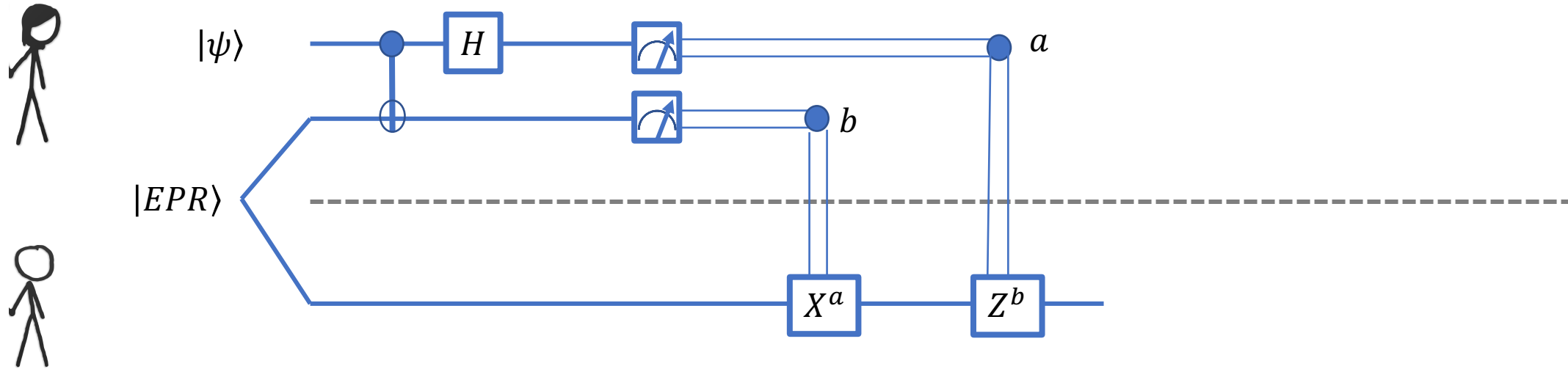
Claim: At the end of protocol, Bob has $|\psi\rangle$.

Quantum Teleportation



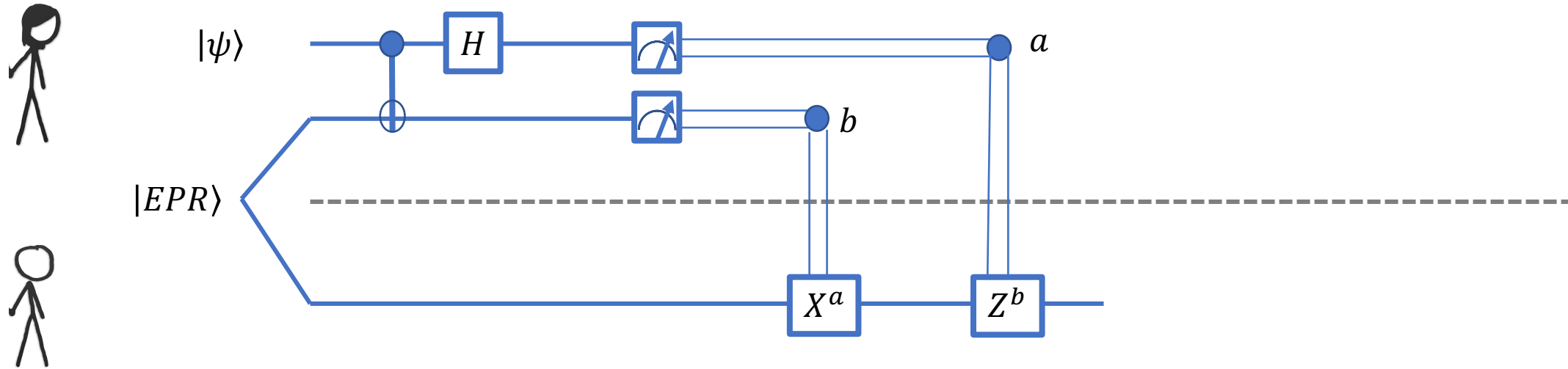
Claim: At the end of protocol, Bob has $|\psi\rangle$.

Quantum Teleportation



Claim: At the end of protocol, Bob has $|\psi\rangle$.

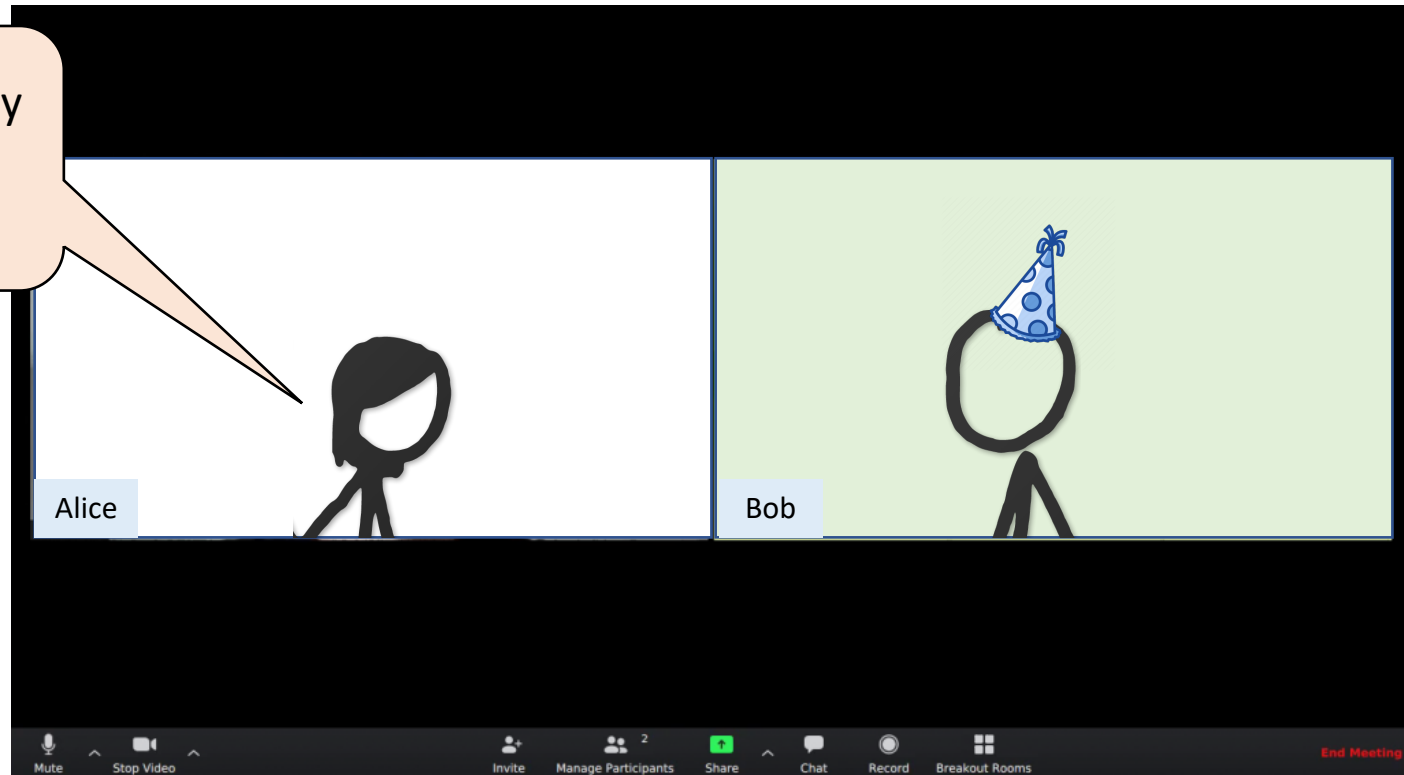
Quantum Teleportation



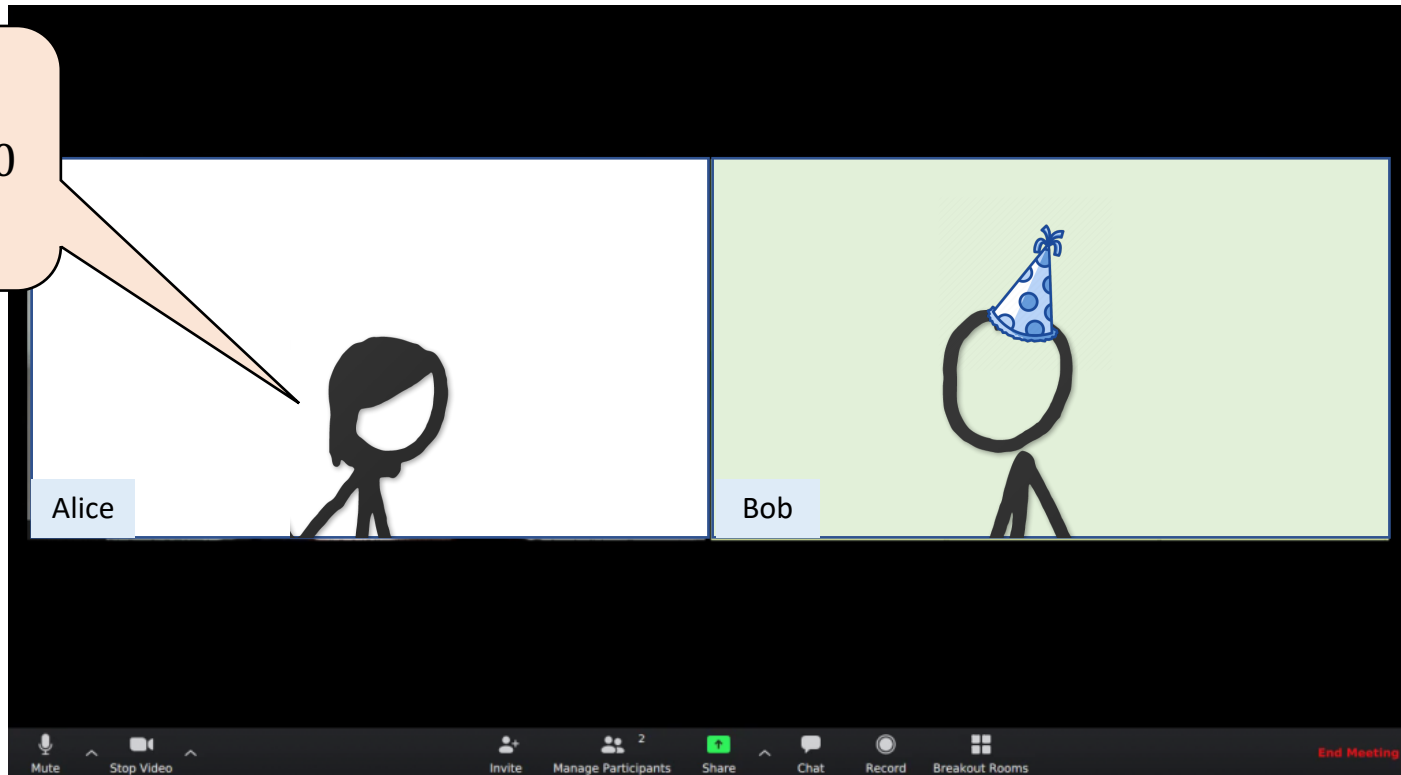
Claim: At the end of protocol, Bob has $|\psi\rangle$.

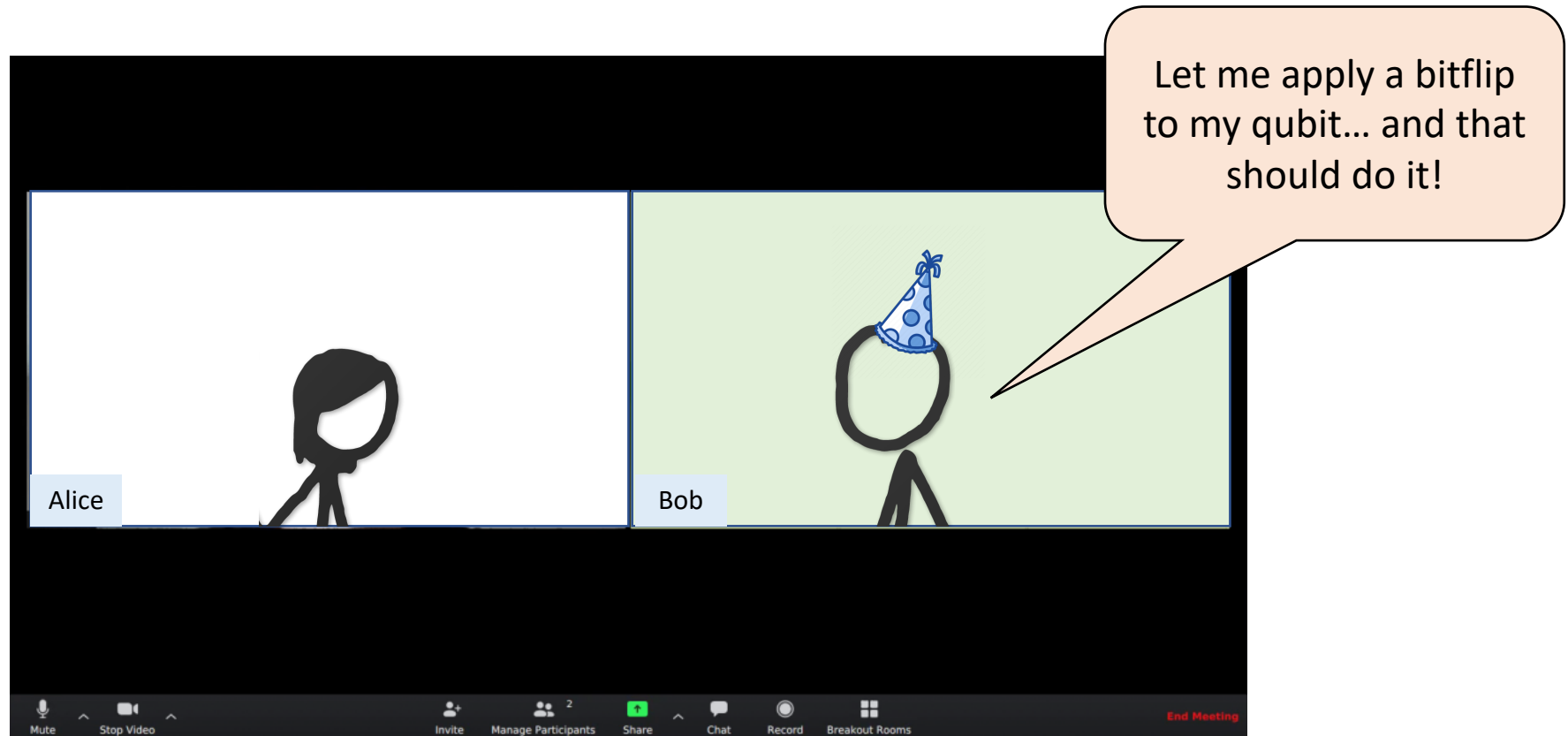
Quantum teleportation does not allow Alice to instantaneously send $|\psi\rangle$ to Bob.
Alice needs to communicate classical bits to Bob!

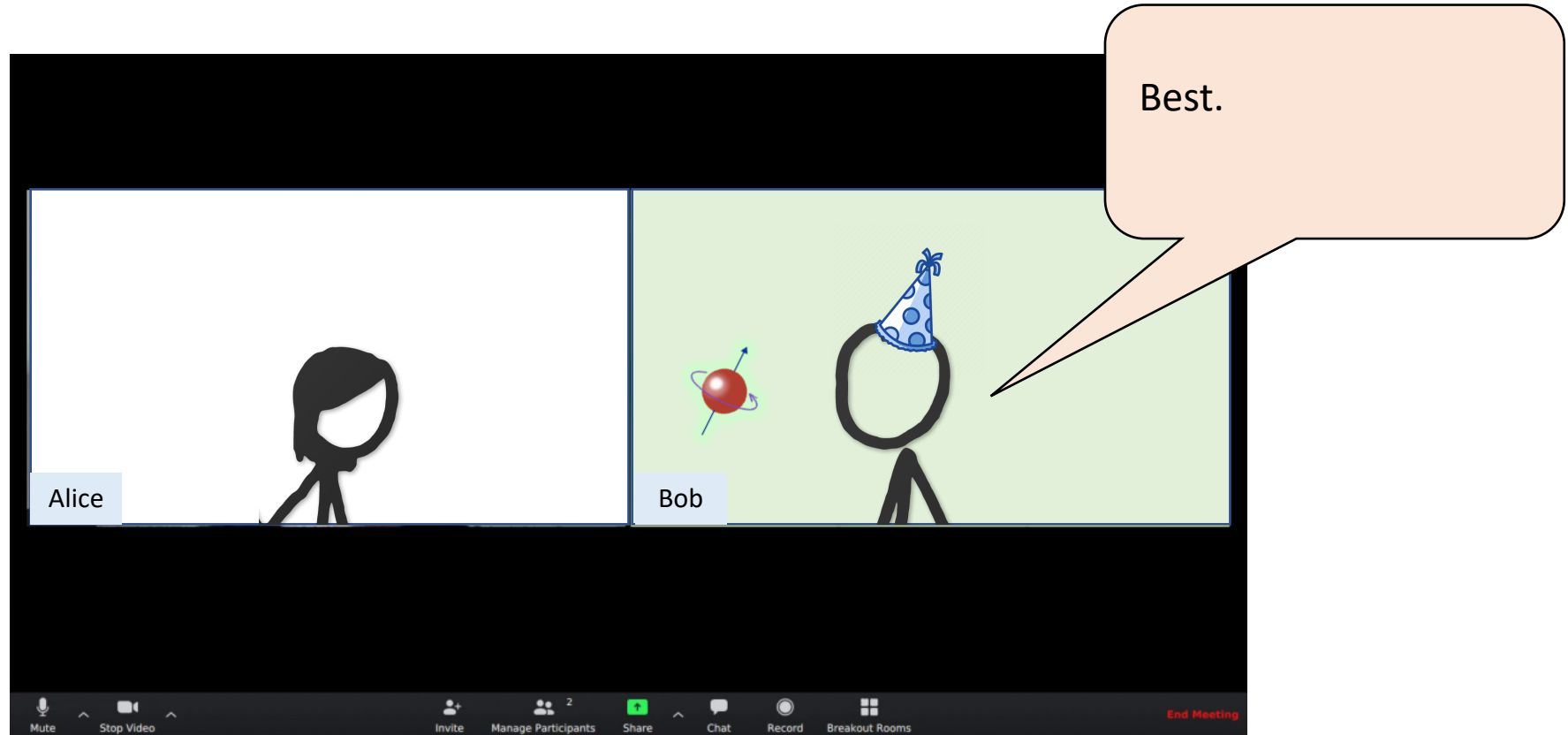
Alrighty... let me apply
the CNOT... then
Hadamard....

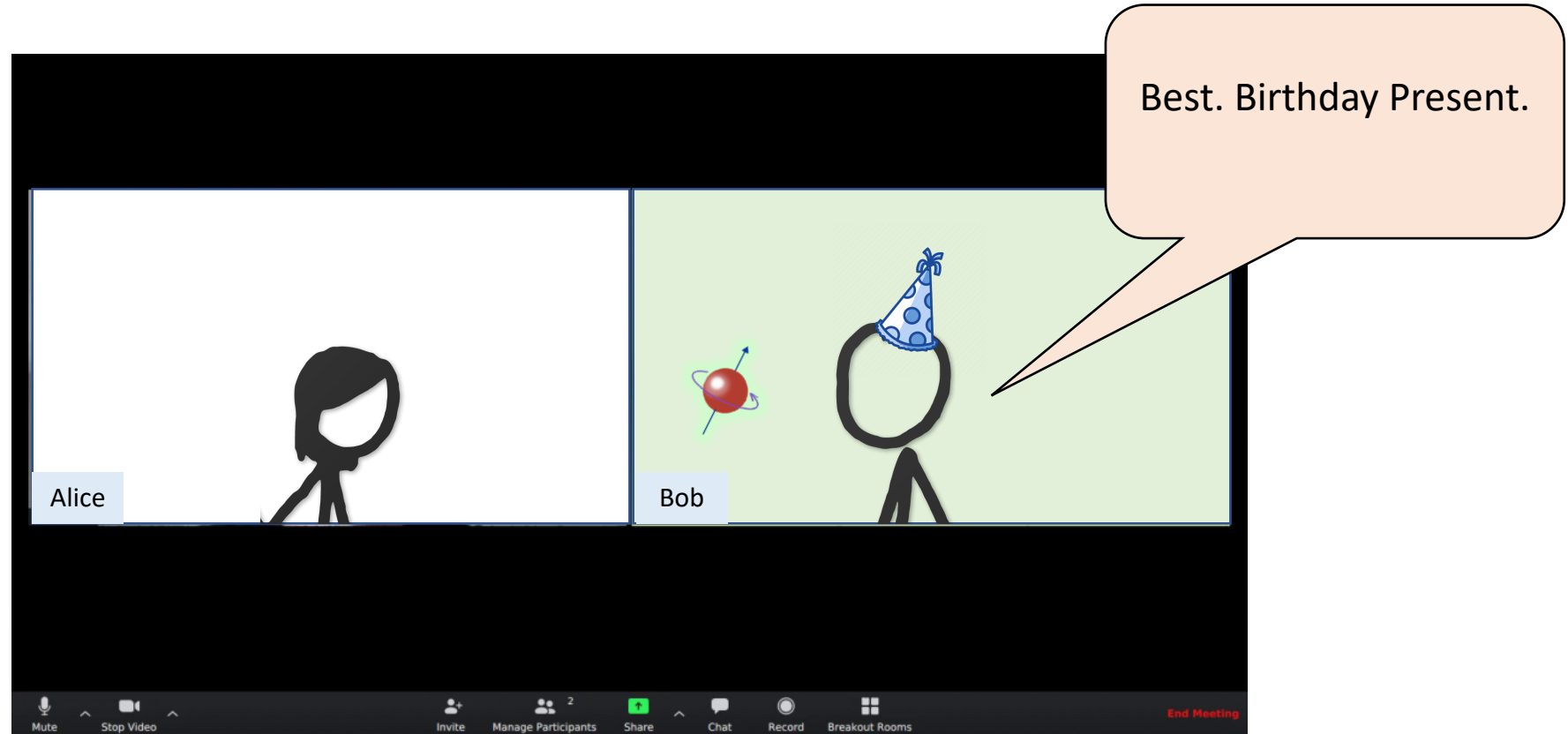


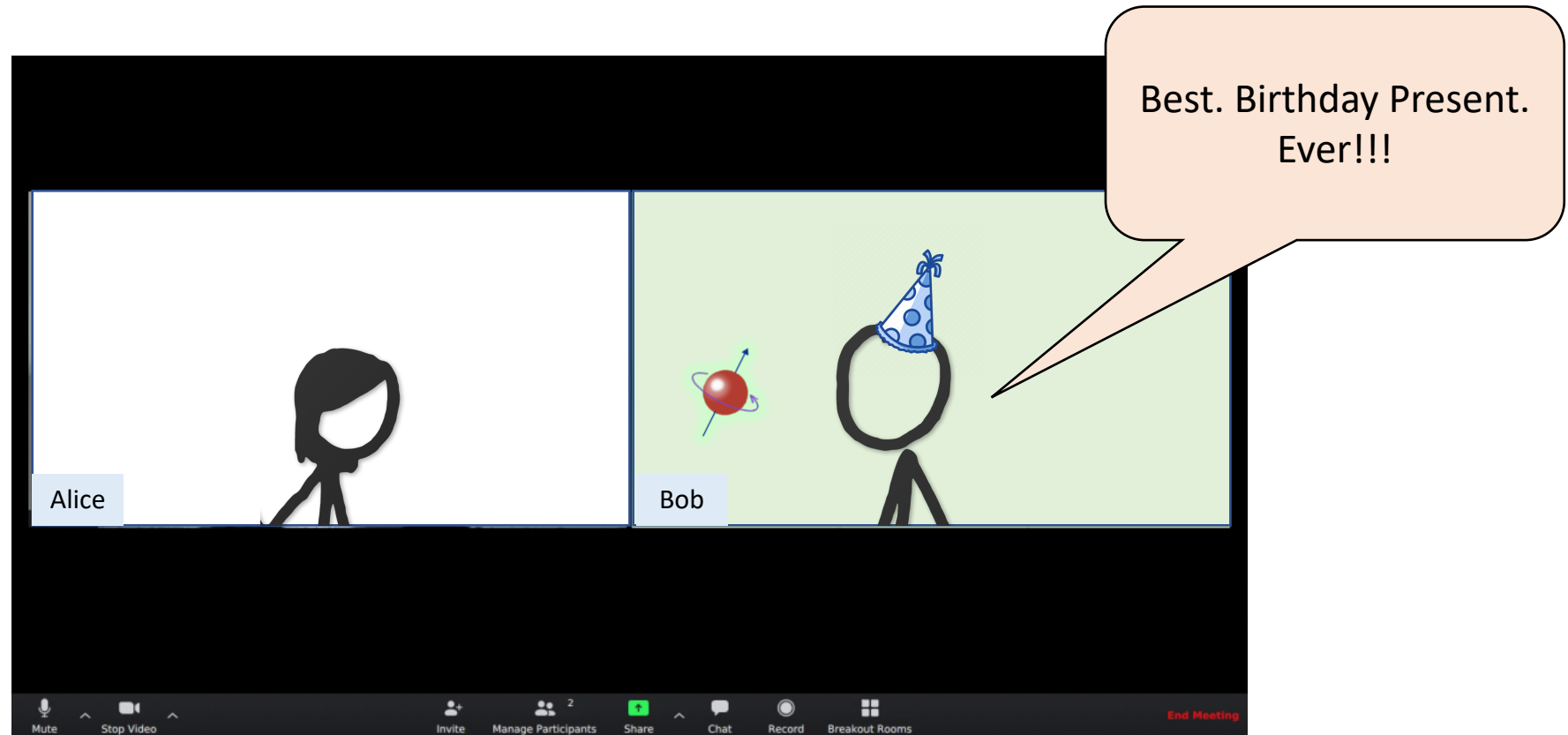
I just measured my
qubits and I got $a = 0$
and $b = 1$.











FIN

Quantum Circuit Model

Quantum gates

- A k qubit-*quantum gate* is a $2^k \times 2^k$ unitary matrix U

- Common single-qubit quantum gates:

- I – identity

- X – bitflip: $|0\rangle \leftrightarrow |1\rangle$

- H – Hadamard: $|0\rangle \mapsto \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$
 $|1\rangle \mapsto \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$

Phase gates

$$Z: |0\rangle \mapsto |0\rangle,$$

$$|1\rangle \mapsto -|1\rangle$$

$$P: |0\rangle \mapsto |0\rangle,$$

$$|1\rangle \mapsto i|1\rangle$$

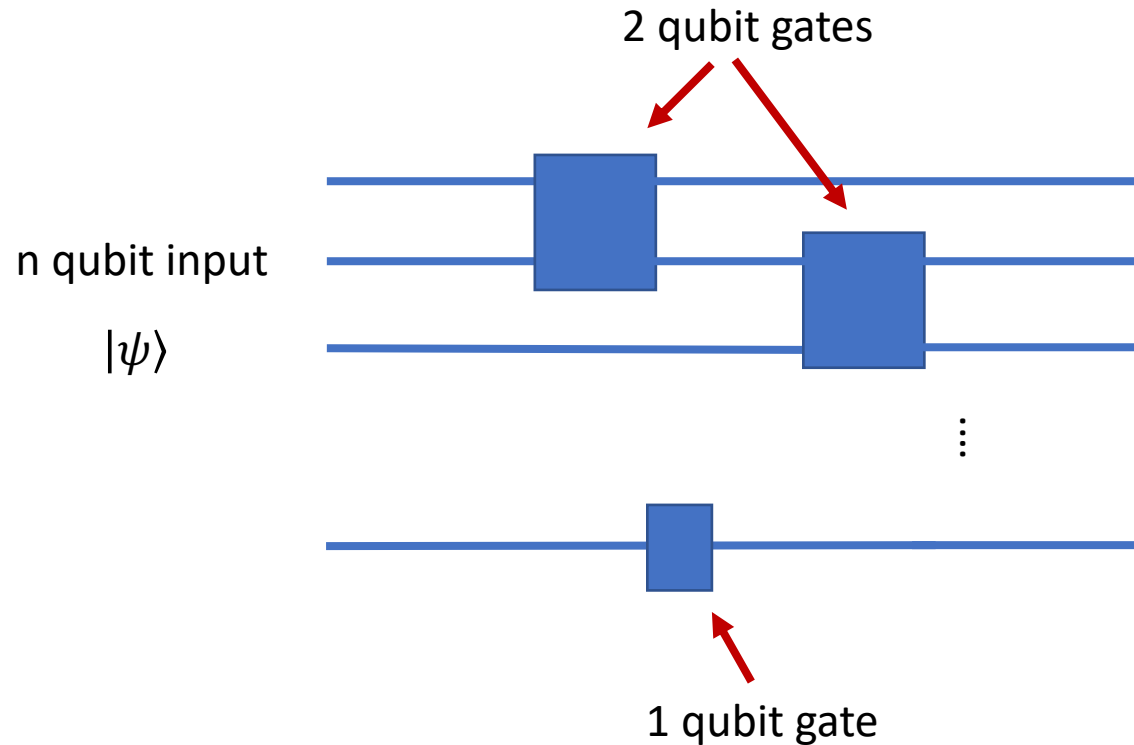
$$T: |0\rangle \mapsto |0\rangle,$$

$$|1\rangle \mapsto e^{\frac{2\pi i}{8}}|1\rangle$$

- Two-qubit gates:

- $CNOT$ – controlled NOT operation: $CNOT|x, a\rangle = |x, a \oplus x\rangle$

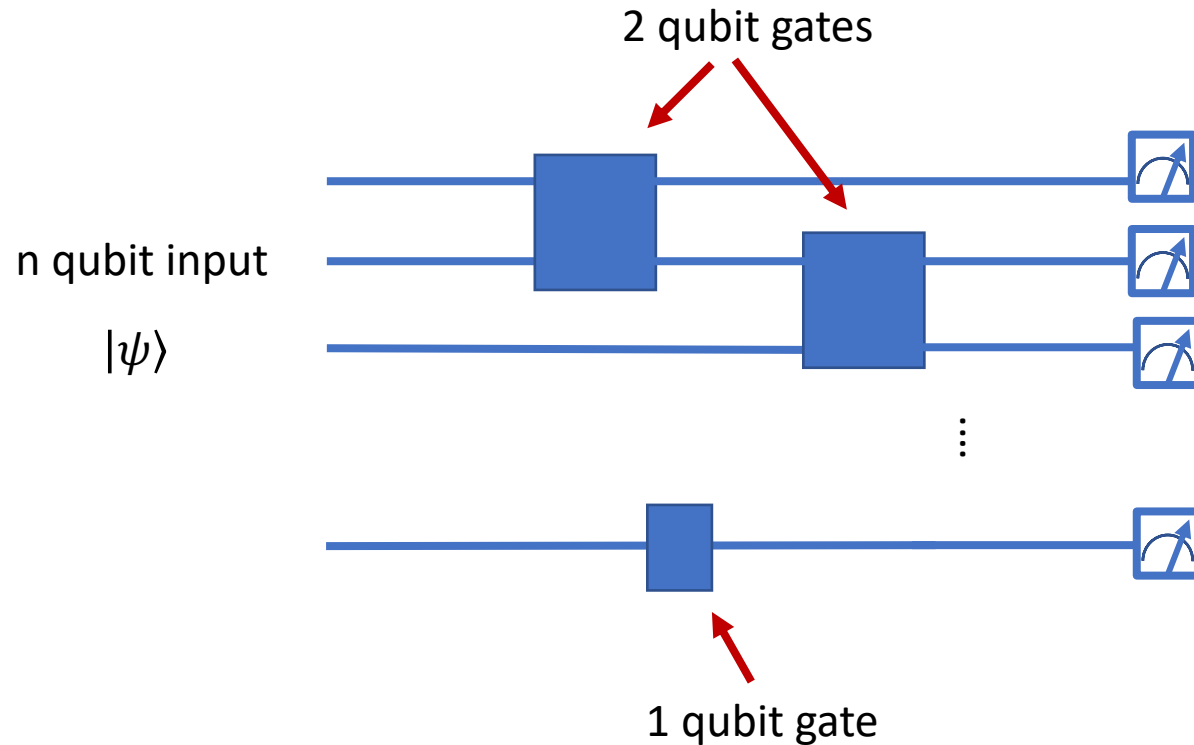
Quantum circuits



- A quantum circuit F consists of an ordered collection of 1- and 2-qubit gates G_1, G_2, \dots applied to subsets of qubits.
- Output of circuit F on input $|\psi\rangle$ is equal to

$$G_m \cdots G_2 G_1 |\psi\rangle$$

Measurements



- At end of computation, if final state is

$$|\varphi\rangle = \sum \beta_x |x\rangle$$

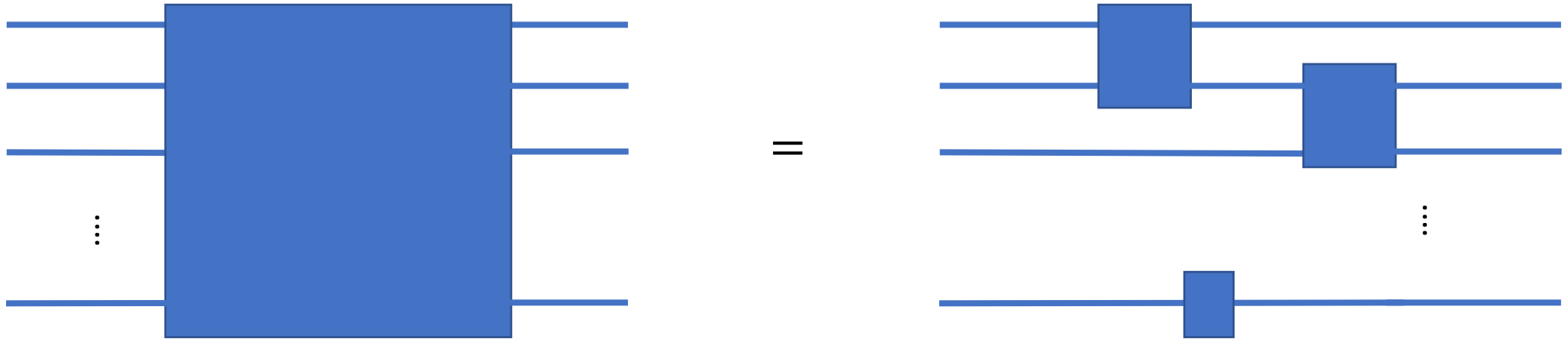
can perform **measurement** to get classical outcome of computation.

- Measurement is probabilistic: obtains outcome $x \in \{0,1\}^n$ with probability $|\beta_x|^2$.
- Measurement is **destructive**: measuring in middle of quantum computation will disturb the state.

We can also allow intermediate measurements (like in quantum teleportation), but for now let's assume that measurements happen at the very end.

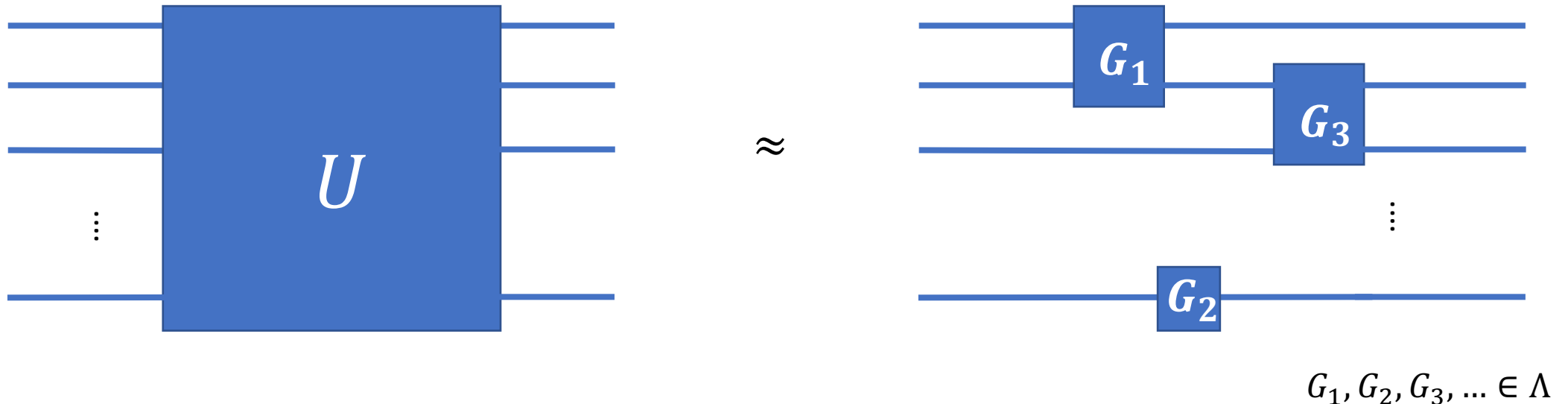
Universal and non-universal gate sets

- Every n -qubit unitary U can be implemented as a quantum circuit consisting of single-qubit gates and CNOT.
- In worst case, such a circuit requires $\approx 4^n$ gates.
- Can use arbitrary single-qubit gates $G \in \mathbb{C}^{2 \times 2}$.



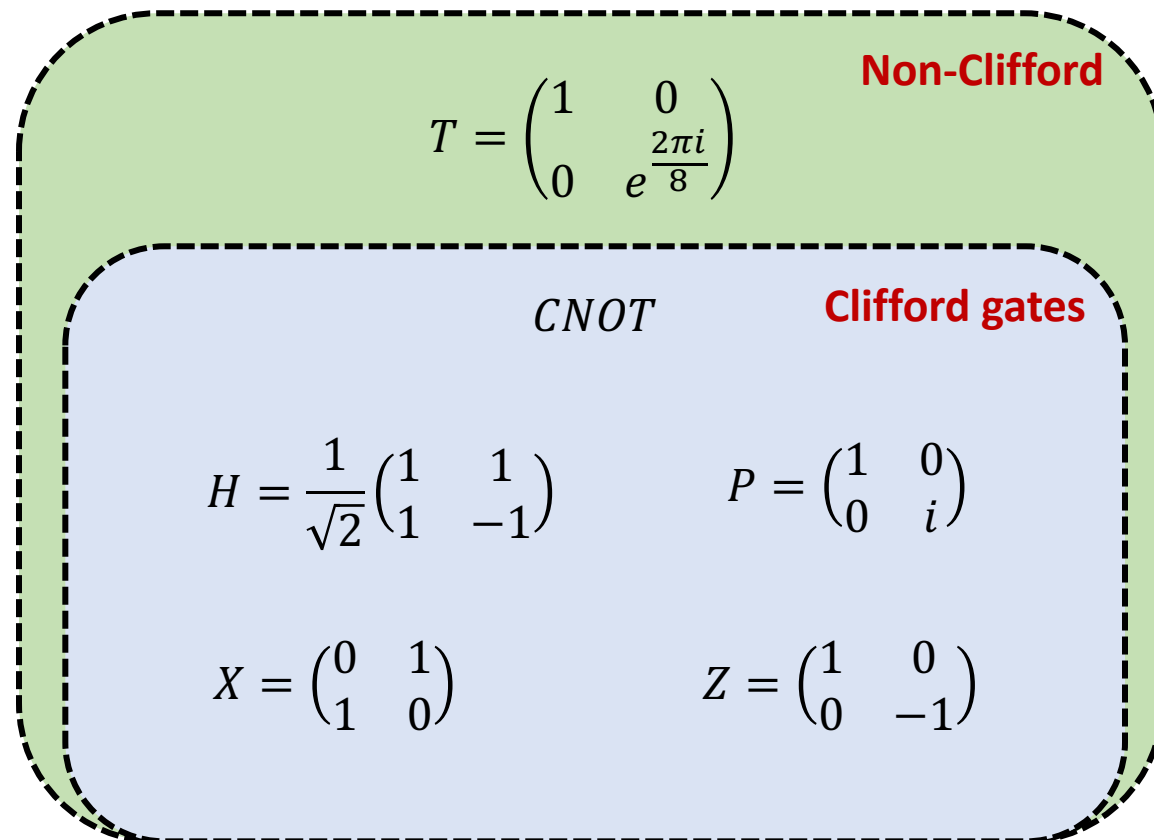
Universal and non-universal gate sets

- In practice, we can only use gates from a fixed, finite set (depending on your hardware).
- A set Λ of gates is **universal** if any unitary (on any number of qubits) can be approximated arbitrarily well by a circuit consisting of gates from Λ .
- A unitary U ϵ -approximates another unitary V if: $\max_{|\psi\rangle} \|U|\psi\rangle - V|\psi\rangle\| \leq \epsilon$



Universal and non-universal gate sets

- **Ex:** $\Lambda = \text{Clifford} \cup \{T\}$ is a universal gate set!
 - Clifford = gates generated by $\{H, P, CNOT\}$



Universal and non-universal gate sets

- **Ex:** $\Lambda = \text{Clifford}$ is *not* universal gate set.
 - Clifford = gates generated by $\{H, P, CNOT\}$

Fact #0: Clifford circuits are not even universal for *classical* computation.

Fact #1: Clifford circuits (with all zeroes input) can be efficiently simulated on classical computers (Gottesman-Knill Theorem).

CNOT

Clifford gates

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

$$P = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}$$

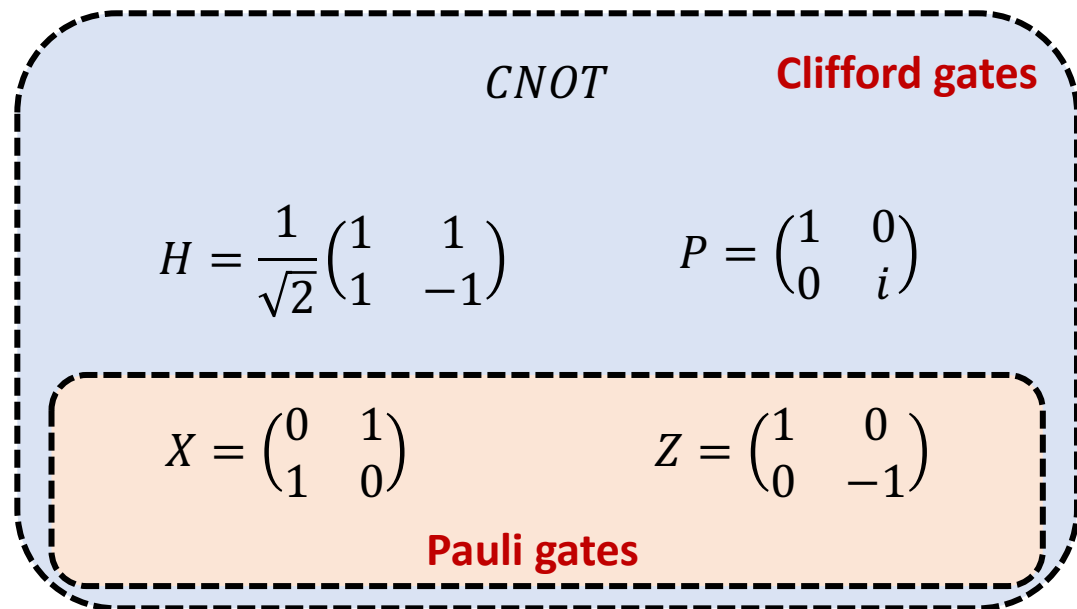
$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Universal and non-universal gate sets

- **Ex:** $\Lambda = \text{Clifford}$ is *not* universal gate set.
 - Clifford = gates generated by $\{H, P, CNOT\}$
- Pauli = gates generated by $\{X, Z\} \subseteq \text{Clifford}$
- n -qubit Pauli unitaries: tensor products of $\{I, X, Y, Z\}$

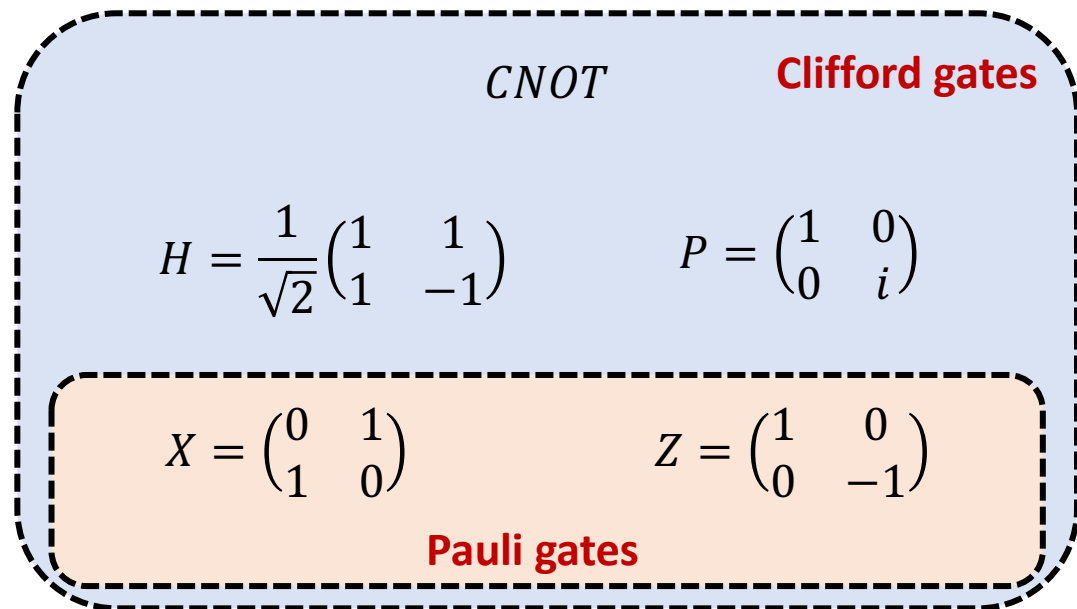
Fact #2: Clifford circuits/unitaries are equivalently defined in terms of their behavior on **Pauli matrices**.



$X \otimes X$
 $X \otimes Z \otimes I \otimes \dots$ } Pauli

Universal and non-universal gate sets

- **Ex:** $\Lambda = \text{Clifford}$ is *not* universal gate set.
 - Clifford = gates generated by $\{H, P, CNOT\}$
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Fact #2: Clifford circuits/unitaries are equivalently defined in terms of their behavior on ***Pauli matrices***.

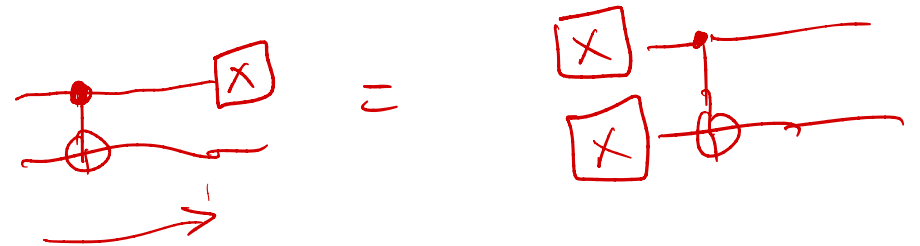
For all Pauli unitaries $\underline{W} = \underline{W_1} \otimes \underline{W_2} \otimes \cdots \otimes \underline{W_n}$

for all n -qubit Clifford unitaries C , there exists another Pauli unitary $\underline{W'} = \underline{W'_1} \otimes \underline{W'_2} \otimes \cdots \otimes \underline{W'_n}$ such that

$$\underline{WC} = \underline{CW'}$$

Ex:

$$XH = HZ$$



Computing classical functions, quantumly

How to compute $f: \{0,1\}^n \rightarrow \{0,1\}^m$ using a quantum circuit?

Can call classical functions as a subroutine using **classical oracles**: define the unitary U_f on $n + m$ qubits: for all $x \in \{0,1\}^n, b \in \{0,1\}^m$,

$$U_f |x, b\rangle = |x, b \oplus f(x)\rangle$$

bitwise XOR. } is a unitary

Ex: $f = \text{AND}, f = \text{NOT}$

$$U_{\text{AND}} |a, b, c\rangle = |a, b, c \oplus ab\rangle.$$

$$U_{\text{NOT}} = \text{CNOT}.$$

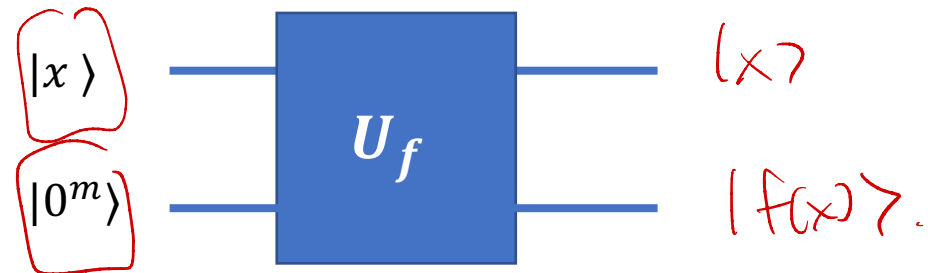
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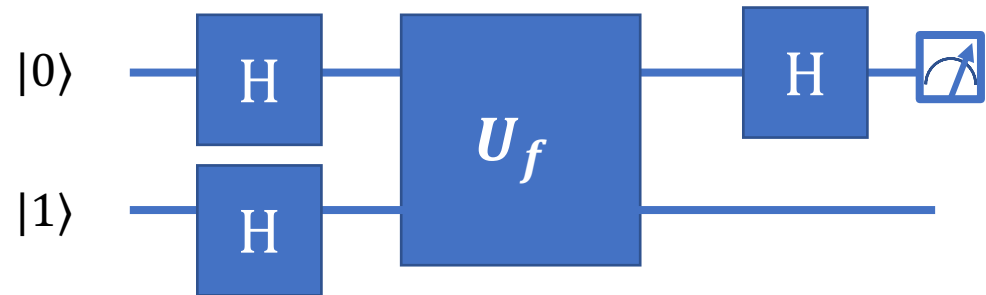
Size s classical circuit computing $f \implies$ There is a size $O(s)$ quantum circuit computing U_f .

Computing classical functions, quantumly

Magic starts happening when classical oracles are queried on a ***superposition*** of inputs.

Deutsch's Problem: Given $f: \{0,1\} \rightarrow \{0,1\}$, determine using one quantum query to U_f whether

- YES case: $f(0) \neq f(1)$
- NO case: $f(0) = f(1)$



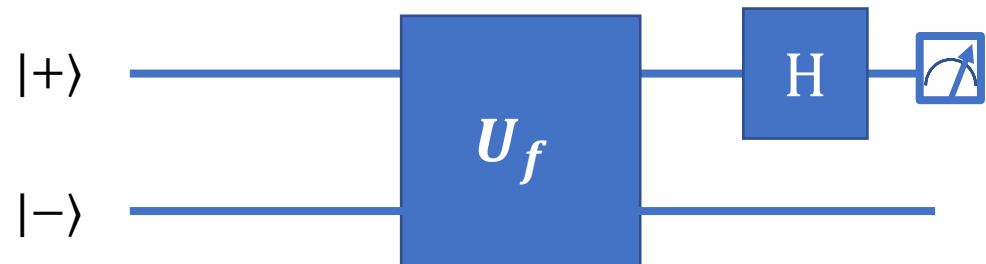
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circuit equivalent to



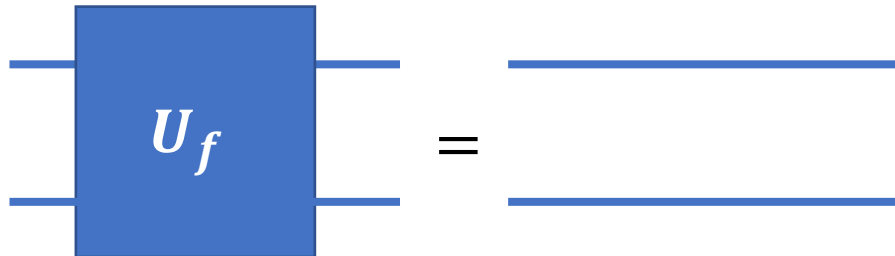
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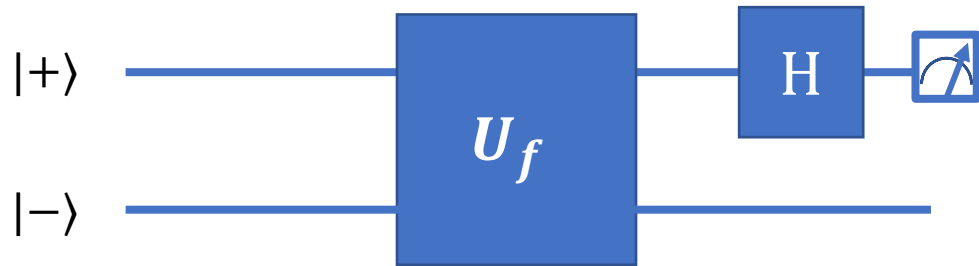
Case 1: $f(0) = f(1) = 0$



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Computing classical functions, quantumly



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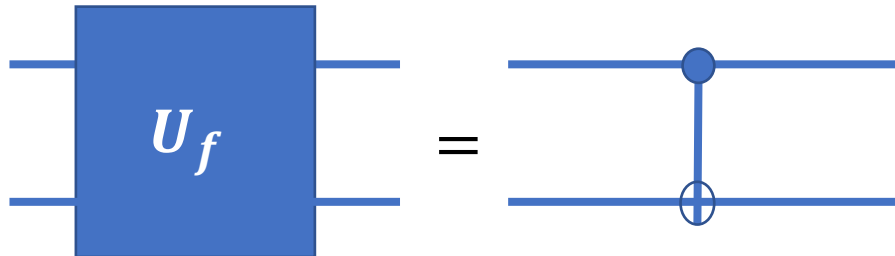
Computing classical functions, quantumly

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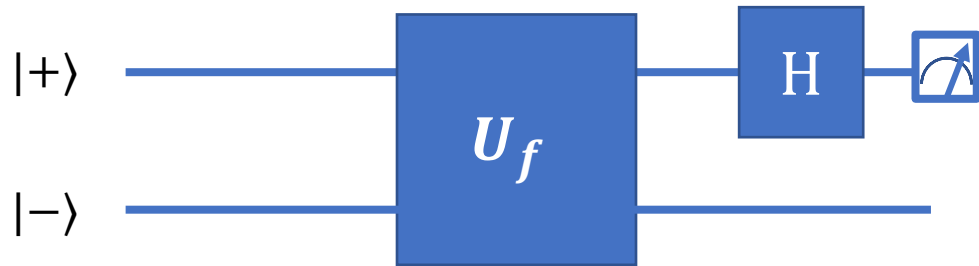
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- NO case: $f(0) = f(1)$

Case 2: $f(0) = 0, f(1) = 1$

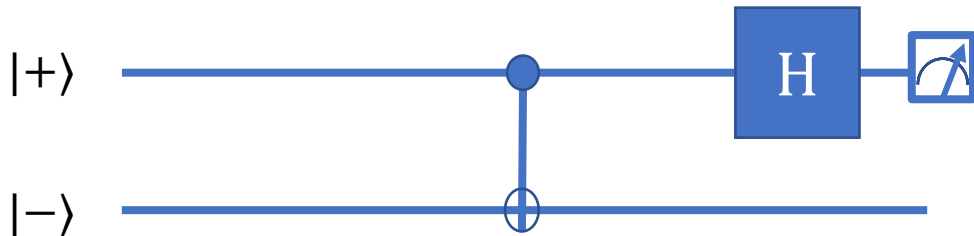


Computing classical functions, quantumly



Case 2: $f(0) = 0, f(1) = 1$

circuit equivalent to



Grover Search

Unstructured search

Search problem: Given black-box access to $f: \{0,1\}^n \rightarrow \{0,1\}$, find x such that $f(x) = 1$.

Classical query complexity: $\Omega(2^n)$

Quantum query complexity: $O(\sqrt{2^n})$

Unstructured search

Search problem: Given black-box access to $f: \{0,1\}^n \rightarrow \{0,1\}$, find x such that $f(x) = 1$.

For boolean functions, we can use different oracle (called *phase* oracle): for all $x \in \{0,1\}^n$

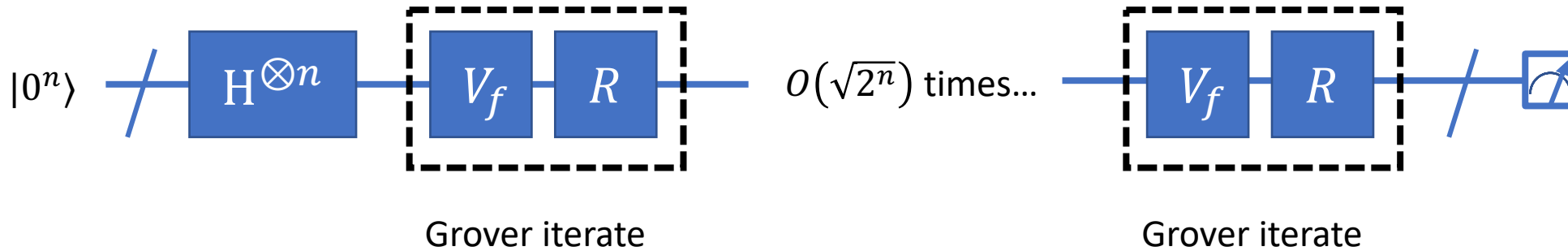
$$V_f|x\rangle = (-1)^{f(x)}|x\rangle$$

XOR oracles and phase oracles are equivalent!

Unstructured search

Search problem: Given black-box access to $f: \{0,1\}^n \rightarrow \{0,1\}$, find x such that $f(x) = 1$.

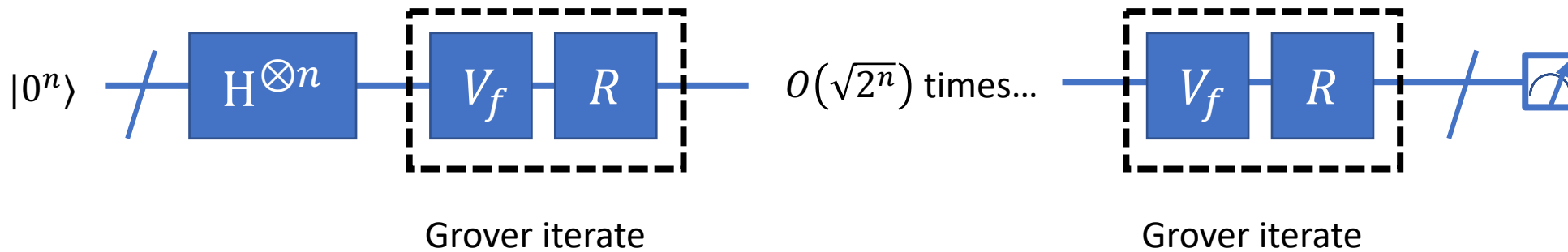
Assume there exists a unique x^* such that $f(x^*) = 1$.



Unstructured search

Search problem: Given black-box access to $f: \{0,1\}^n \rightarrow \{0,1\}$, find x such that $f(x) = 1$.

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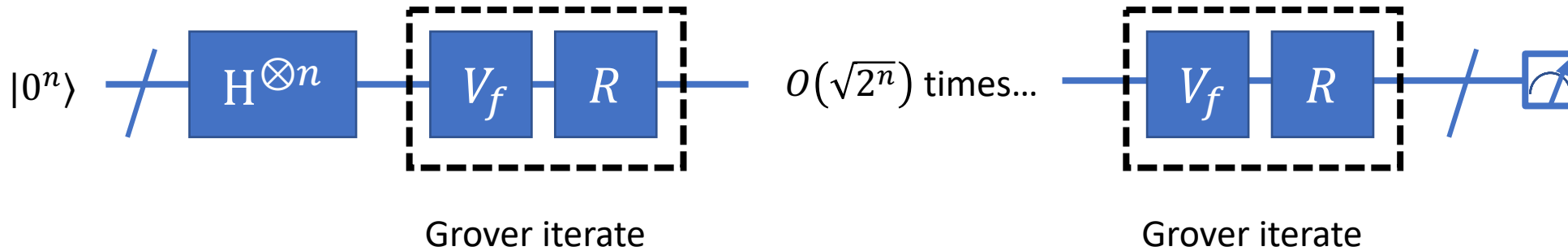


$$|0^n\rangle \xrightarrow{H^{\otimes n}} = |+\rangle^{\otimes n} = \frac{1}{\sqrt{2^n}} \sum_x |x\rangle \quad \} \quad |s\rangle$$

Unstructured search

Search problem: Given black-box access to $f: \{0,1\}^n \rightarrow \{0,1\}$, find x such that $f(x) = 1$.

Assume there exists a unique x^* such that $f(x^*) = 1$.



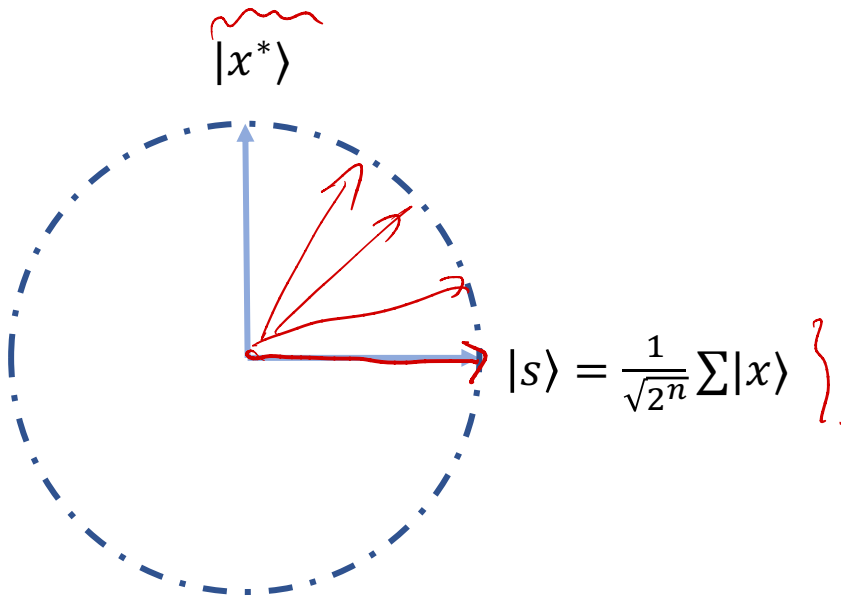
"diffusion operator", "inversion about the mean",...

$$\text{---} \boxed{R} \text{---} = 2|s\rangle\langle s| - I \quad \left. \vphantom{2|s\rangle\langle s| - I} \right\} \text{unitary'}$$

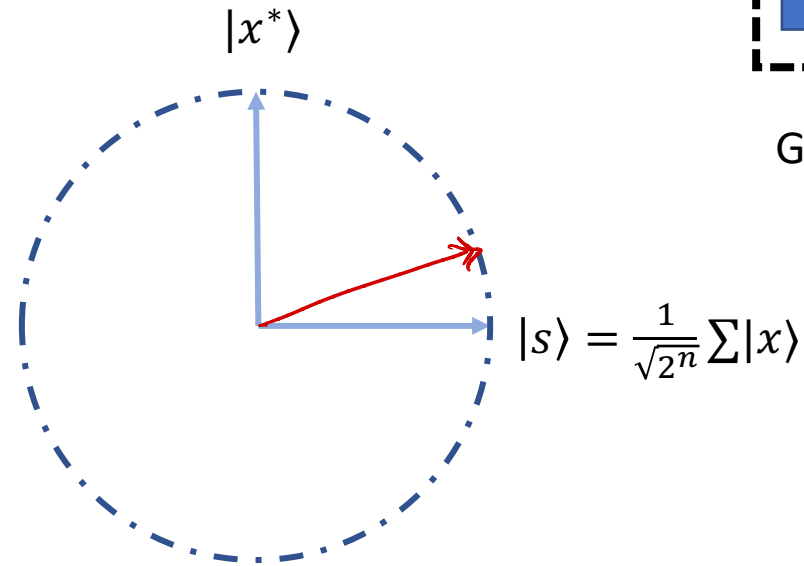
Unstructured search

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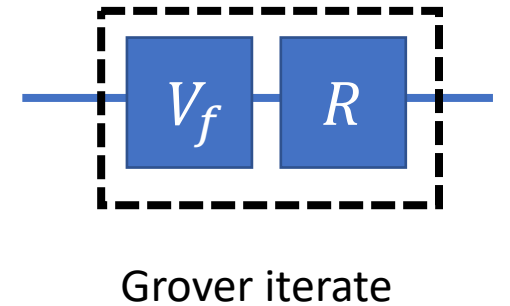
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Starting state of algorithm



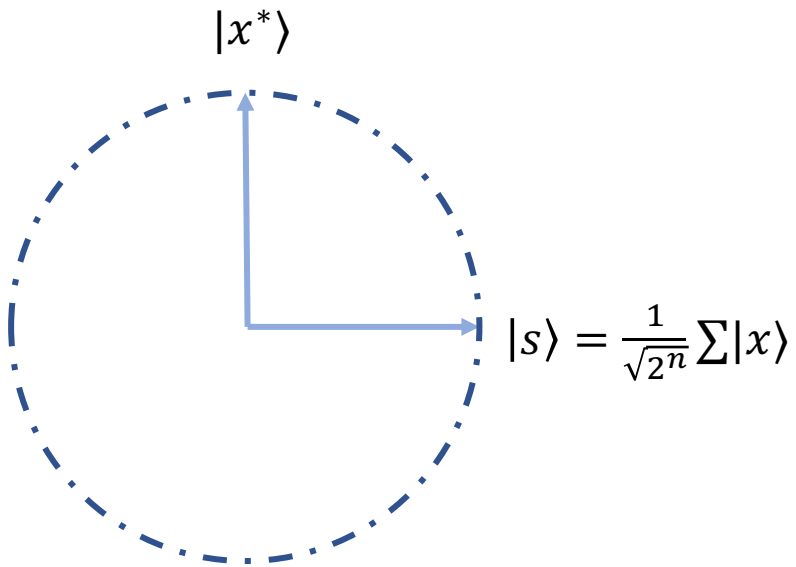
After one Grover iterate



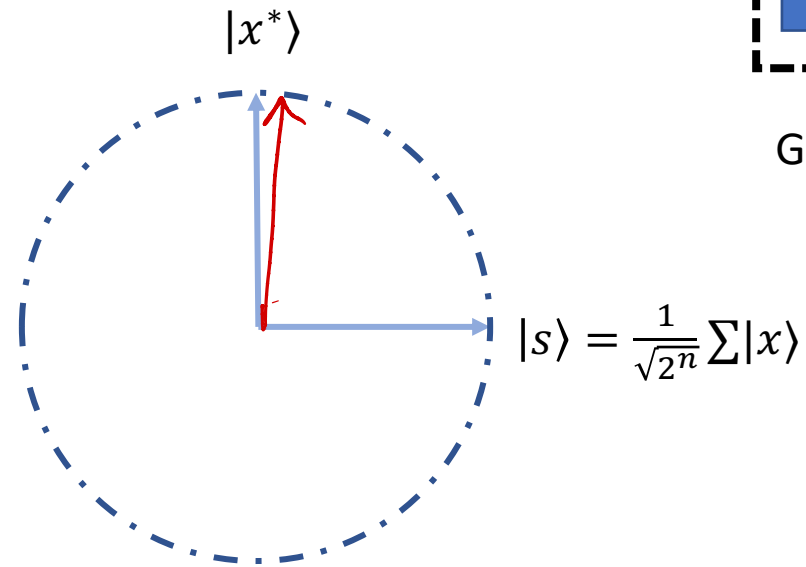
Unstructured search

Search problem: Given black-box access to $f: \{0,1\}^n \rightarrow \{0,1\}$, find x such that $f(x) = 1$.

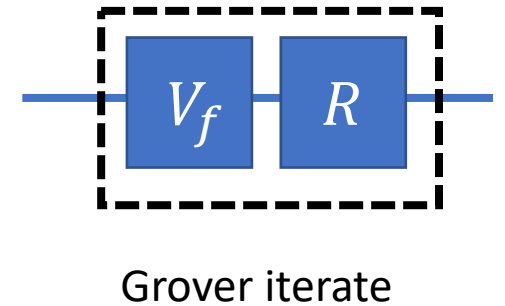
Assume there exists a unique x^* such that $f(x^*) = 1$.



Starting state of algorithm



after $O(\sqrt{2^n})$ iterates



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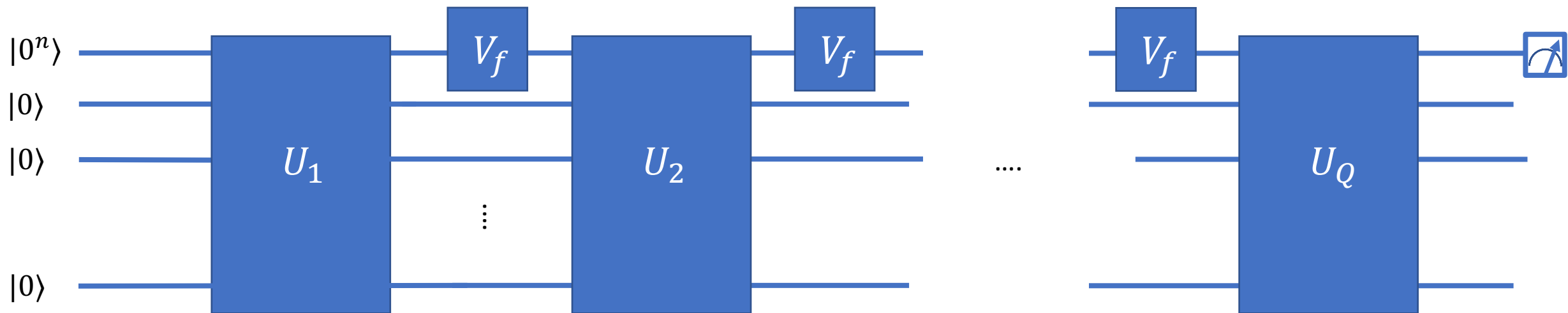
What if there are T solutions x such that $f(x) = 1$?

- If T is known before hand, run $O\left(\sqrt{\frac{2^n}{T}}\right)$ queries.
- If T is unknown, then using more clever algorithm, can still find solution in $O\left(\sqrt{\frac{2^n}{T}}\right)$ queries!

Quantum lower bound for unstructured search

Grover's algorithm is optimal (in terms of query complexity) for solving the unstructured search problem: $\Omega(\sqrt{2^n})$ queries are needed!

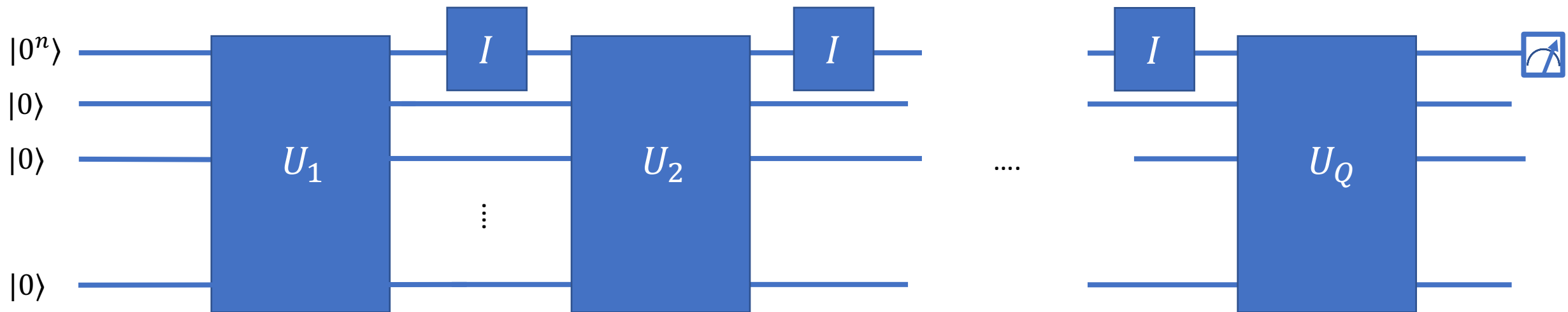
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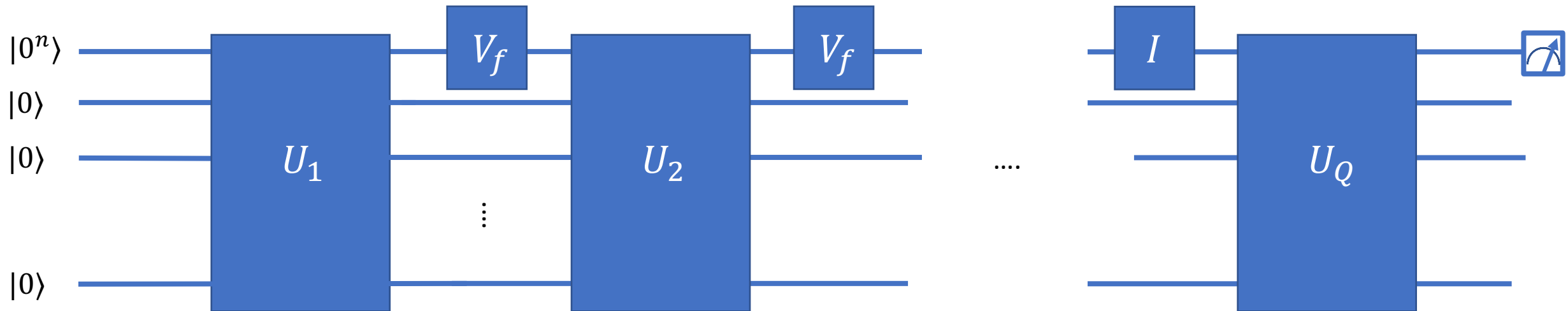
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Generalizations of Grover search

- Quantum Counting
- Amplitude Amplification

